



Lagrangian Studies of the Southern Stratosphere



C. Roberto Mechoso

Dept. Atmospheric and Oceanic Sciences, UCLA

Acknowledgements: Alvaro de la Camara (NCAR, USA), Ana M. Mancho (ICMAT, Spain), Celal Konor (CSU), Ross Heikes (CSU, USA), Anirban Guha (Indian Inst. of Tech.) and Victor Garcia (UAM, Spain)

Outline

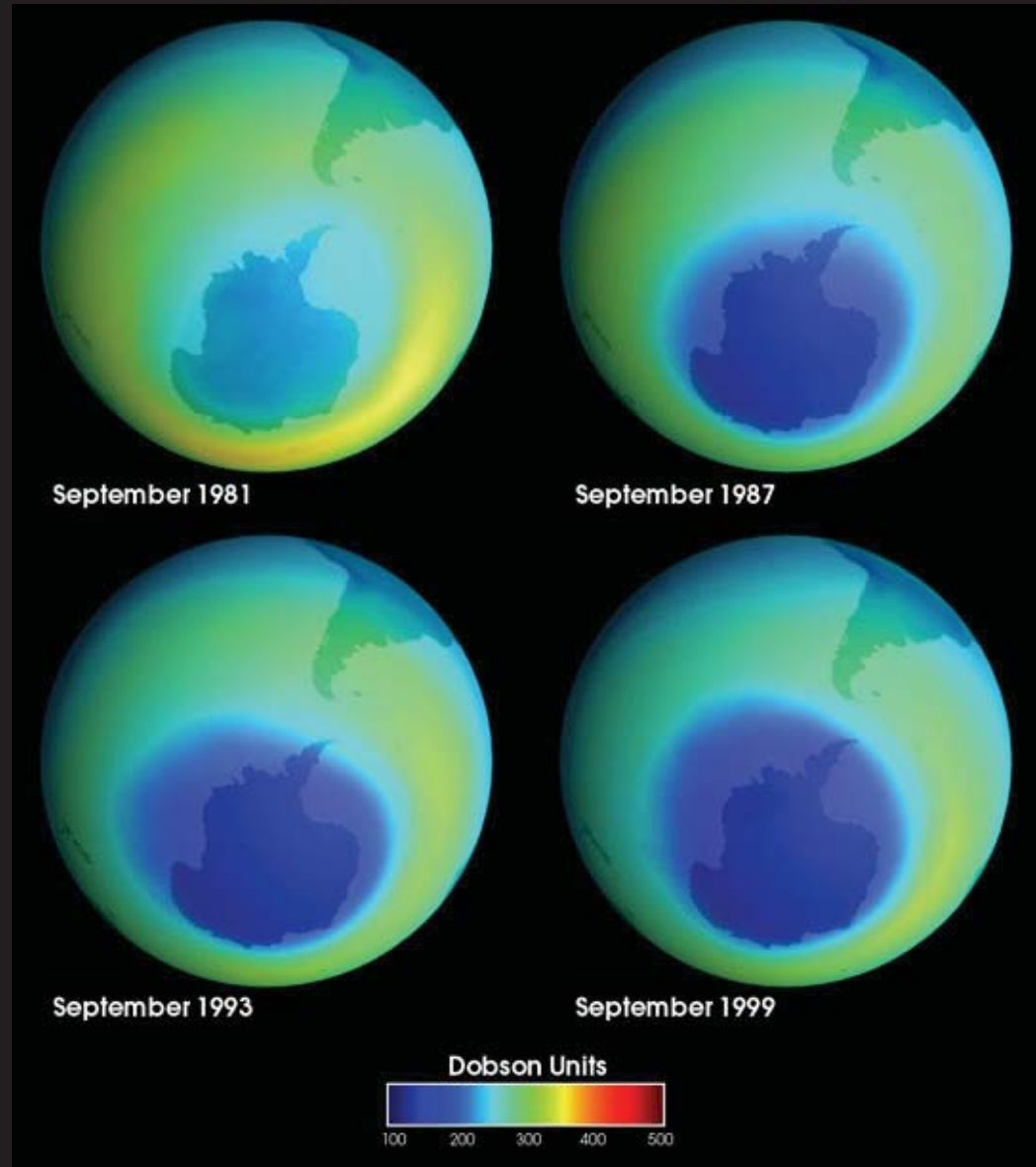
1. The southern spring stratosphere
2. Trajectories inside the polar night vortex
3. Transport across
4. Rossby Wave Breaking
5. Research in progress

The Ozone Hole Over Antarctica

- Why is O₃ destroyed?
- Why in southern Spring?
- Why over Antarctica?

The “hole” seems to be filling up after the ban of pollutants.

*Low ozone levels shown in blue



WINTER POLAR NIGHT JET



**Strong westerlies
(100 m/s)**

**The vortex is
almost circumpolar**

**The pole is dark
and colder than the
equator**

≈30 km high

SUMMER



**Strong easterlies
(50 m/s)**

**The vortex is
almost
circumpolar**

**The pole is
lighted and
warmer than the
equator**

≈30 km high

Typical Flow Regimes in the SH Stratosphere during Spring

September



October



≈ 30 km high

Mechoso et al. (QJRMS; 1988)

SPRING

Vortices are not circumpolar

The pole is connected with warmer air in lower latitudes: “The Final Warming”

The “Antarctic Ozone Hole” forms wherein the cyclonic vortex

≈30 km high

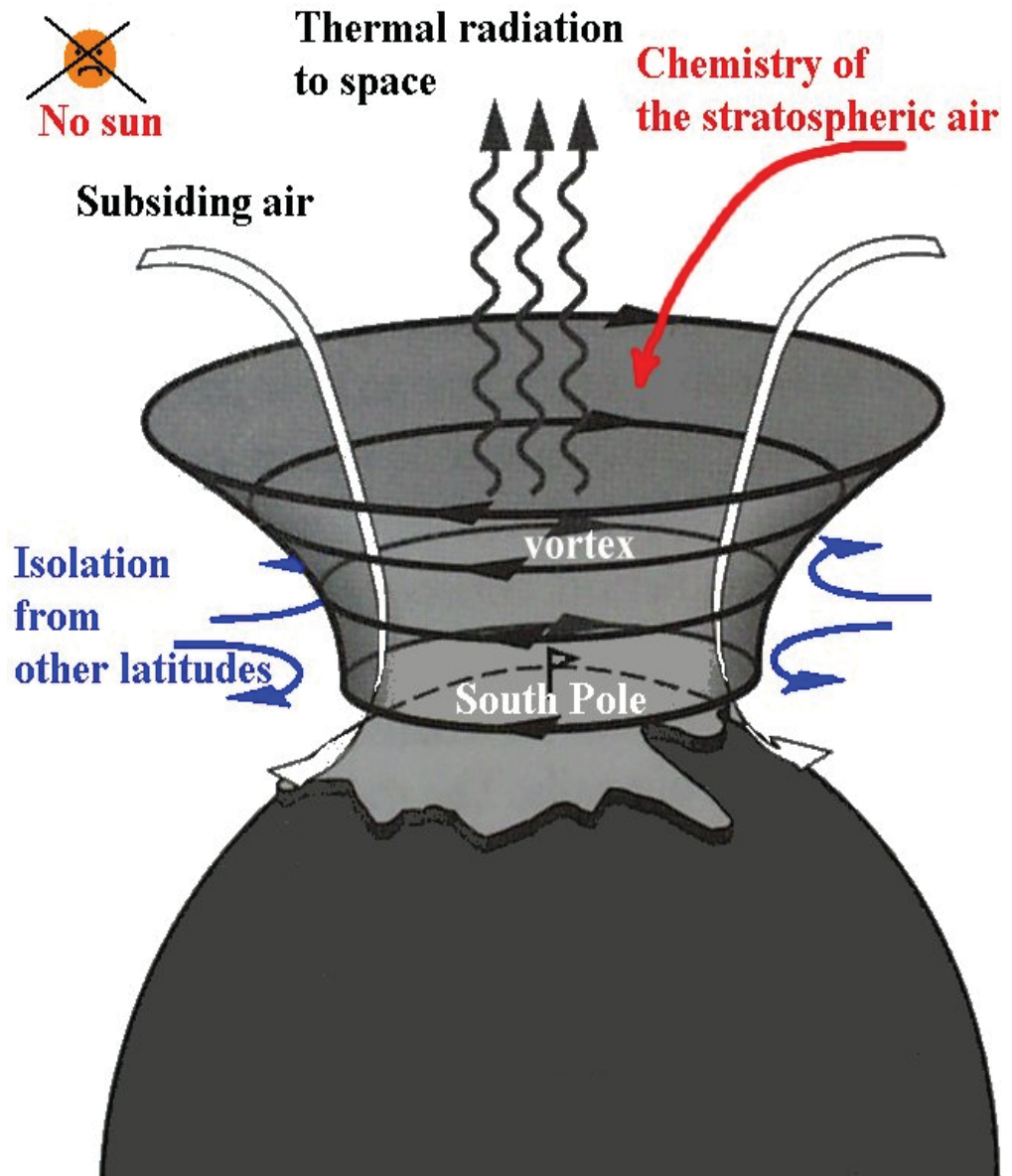


The Antarctic Ozone Hole

The polar night vortex provides a “containment vessel” for the polar air.

Inside this vessel, chemical reactions triggered by the return of the sun in spring destroy the ozone.

How “leaky” is this vessel?



Dynamics: Motion invariants for adiabatic and frictionless flows

Potential Temperature

$$\theta = T \left(p_0 / p \right)^{R/c_p}$$

$$\frac{D\theta}{Dt} = 0$$

Potential Vorticity

$$PV = -g \left(f + \mathbf{k} \cdot \nabla_{\theta} \times \mathbf{v} \right) \left(\partial p / \partial \theta \right)^{-1}$$

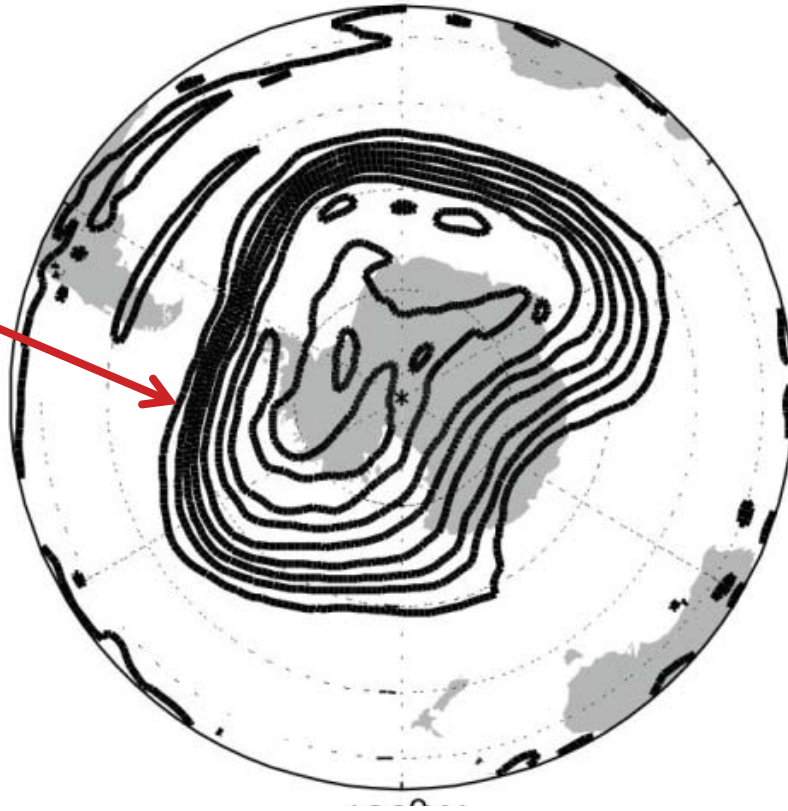
$$\frac{D(PV)}{Dt} = 0$$

These are good approximations in the stratosphere up to ~15 days

Potential Vorticity (PV) at 475 K

The vortex edge is defined as the location of stronger PV gradients

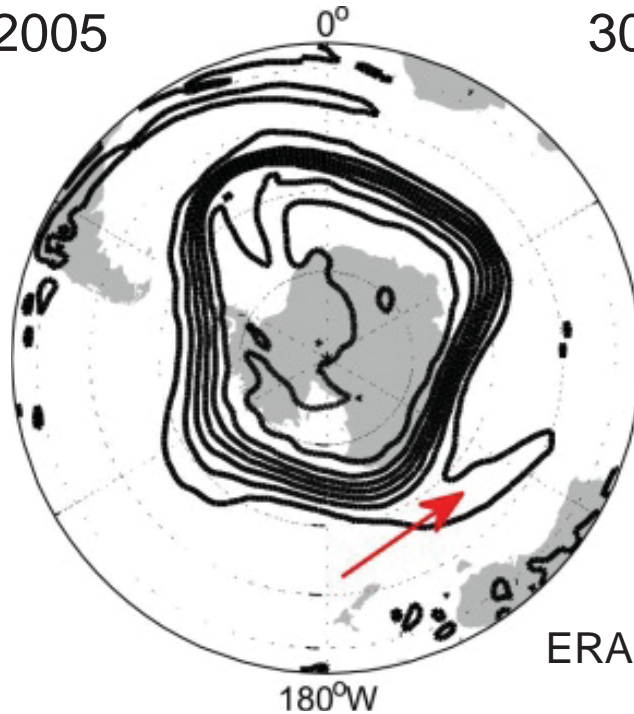
October 15, 2010



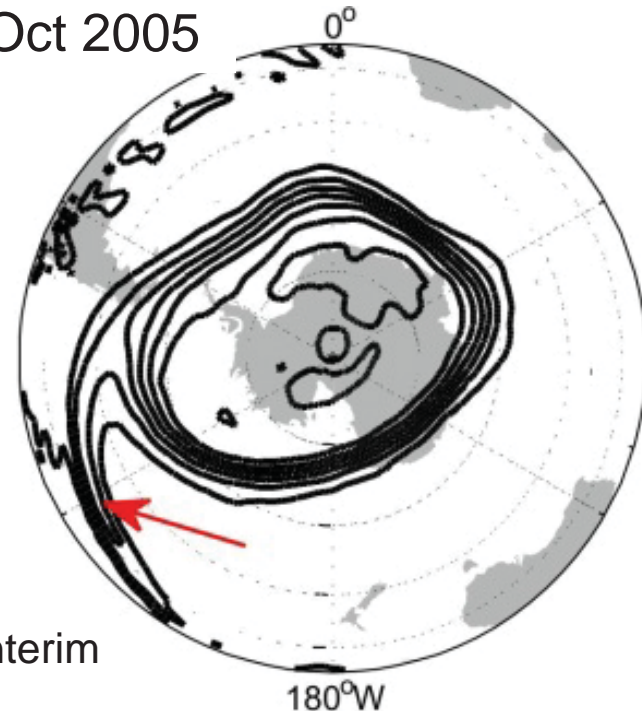
Deviations of the shape of the vortex from a circular shape are due to the presence of planetary-scale (Rossby) waves.

A closer look at PV (475K) reveals finer structures

15 Oct 2005

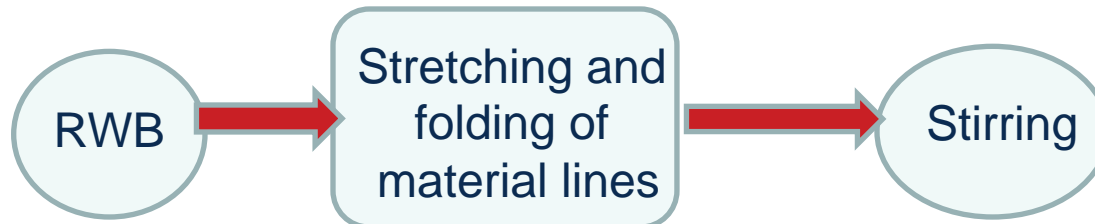


30 Oct 2005



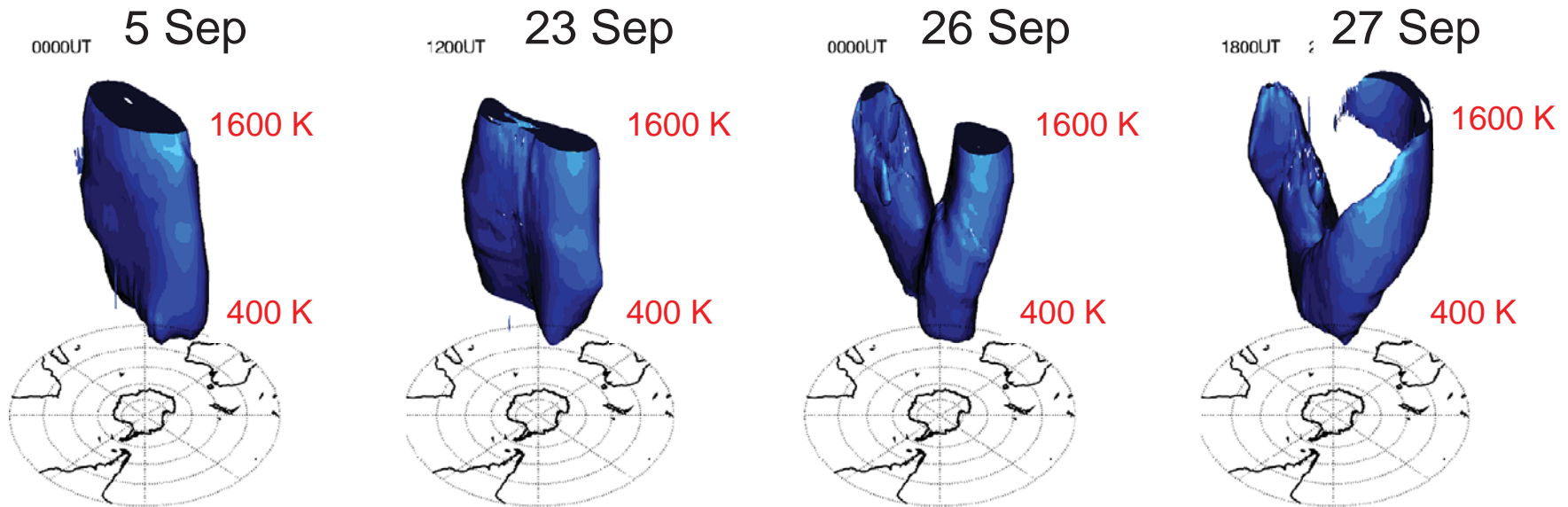
ERA-Interim

Around the vortex there is the surf zone

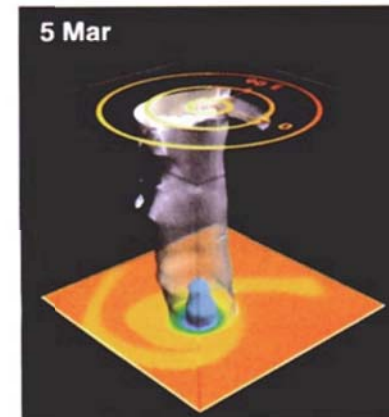
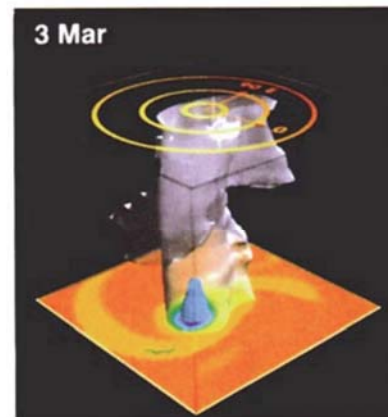
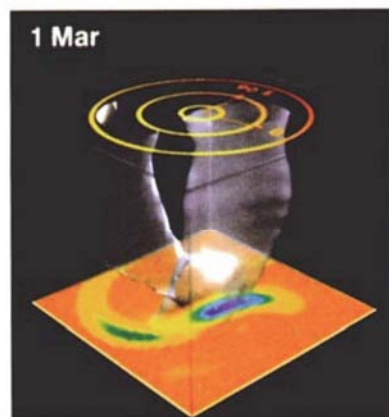
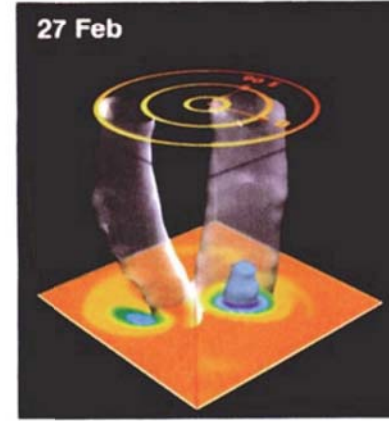
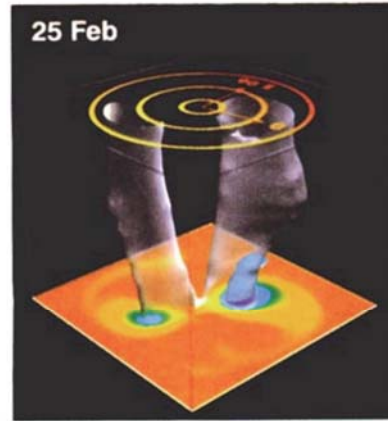
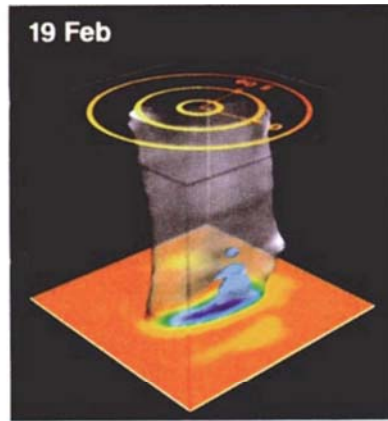
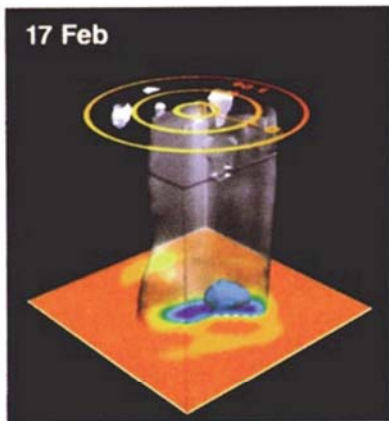


Rossby Wave Breaking

Southern Hemisphere: 3D PV Evolution in September 2002



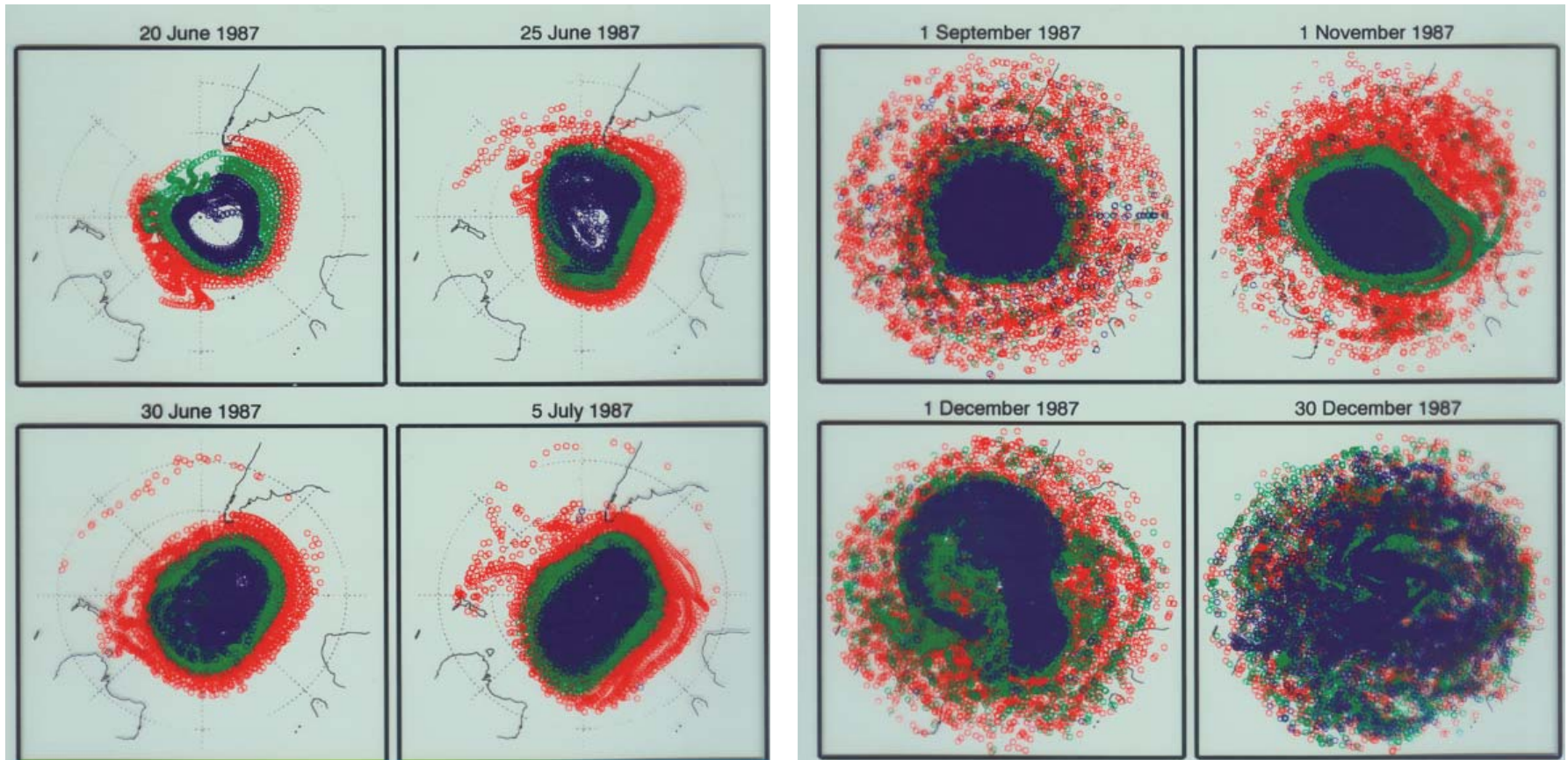
Dataset: ECMWF operational analyses



3D View of the
PNJ splitting
using PV

**Northern
Hemisphere**
February 1979

Simulated “Balloon” Trajectories at 10hPa



“Balloon” means trajectories $\mathbf{x}(t)$ on isopycnal surfaces

$$\frac{d\mathbf{x}(t)}{dt} = \mathbf{v}(t, \mathbf{x})$$

Where the velocity field is provided by a numerical model.

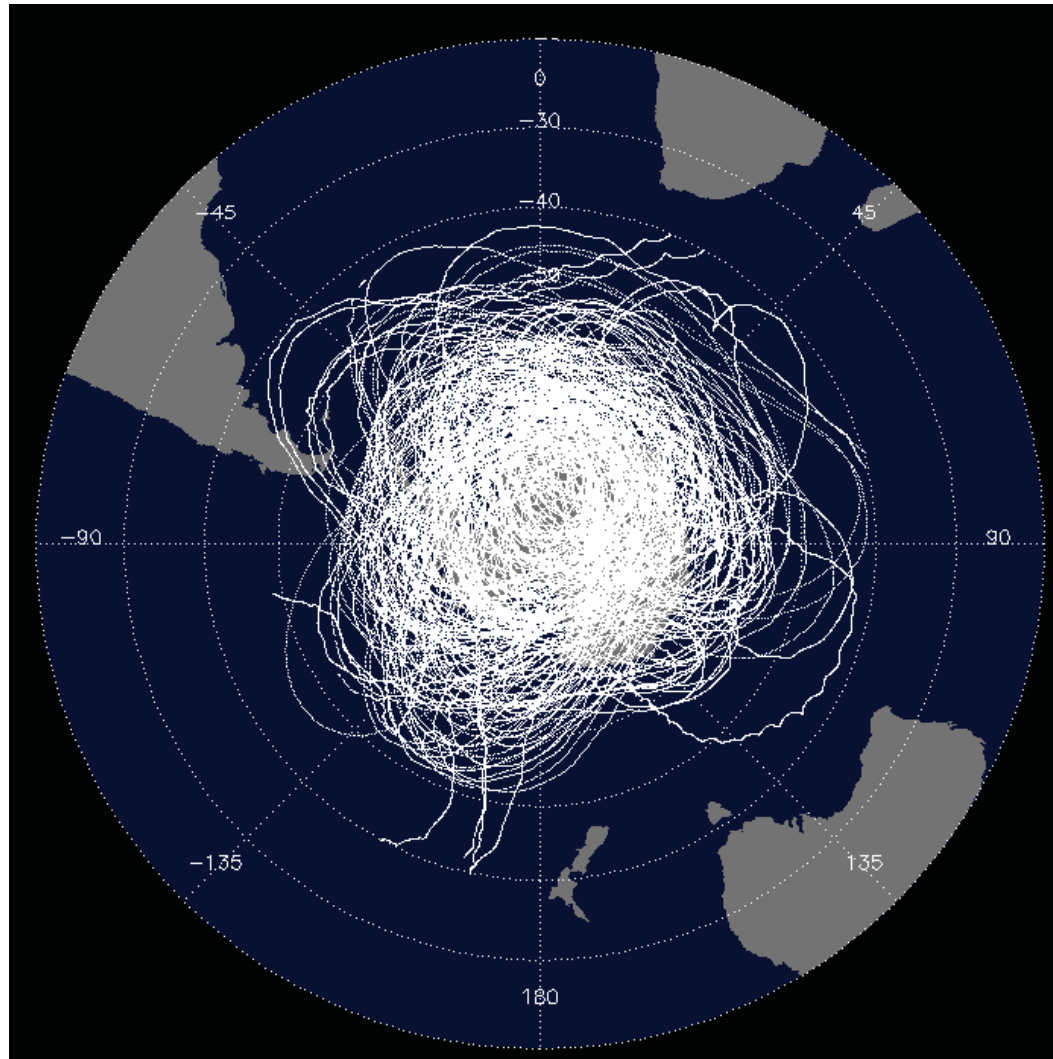
VORCORE and CONCORDIASI field campaigns provided approximations to the behavior of fluid parcels



The campaigns launched 27 and 19 super-pressure balloons from McMurdo, Antarctica, during Sep-Oct 2005 and 2010, respectively.

The balloons drifted on isopycnic surfaces at about 50 Hpa inside the stratospheric polar vortex, gathering meteorological information (temperature, pressure) and position by GPS.

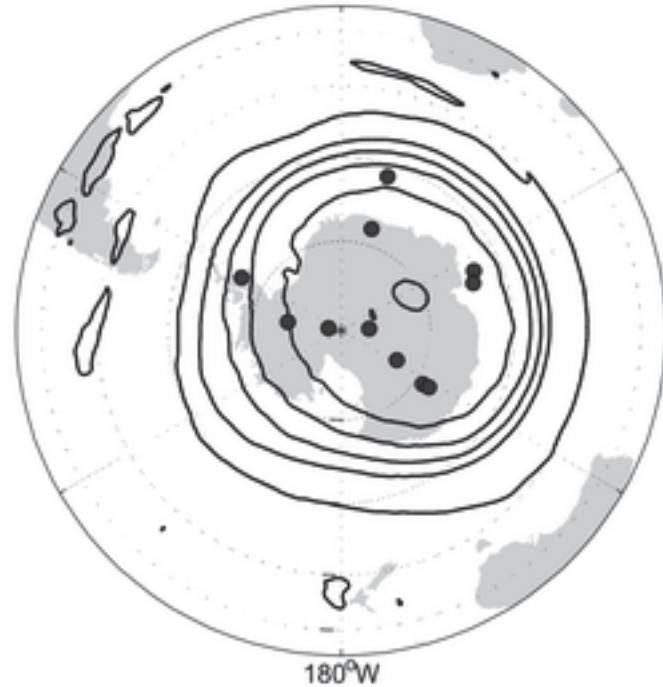
Trajectories of VORCORE balloons



Balloons and PV at 475K



Vorcore
28/Oct/2005

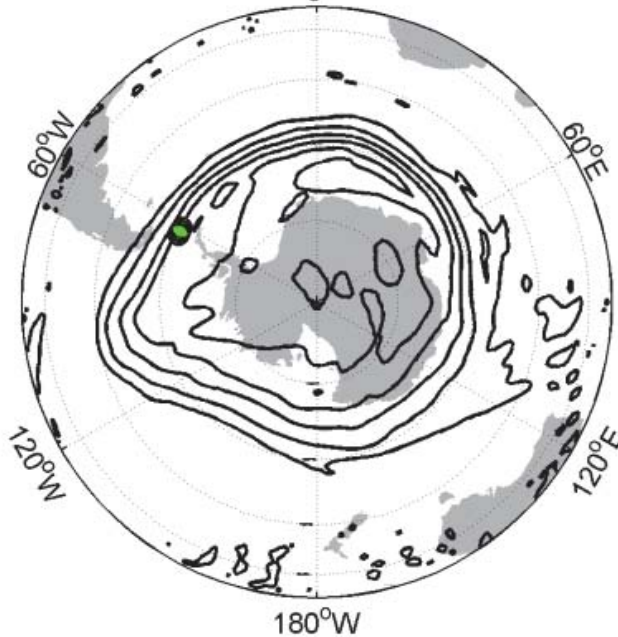


Concordiasi
26/Oct/2010

How can the air parcels cross the vortex?

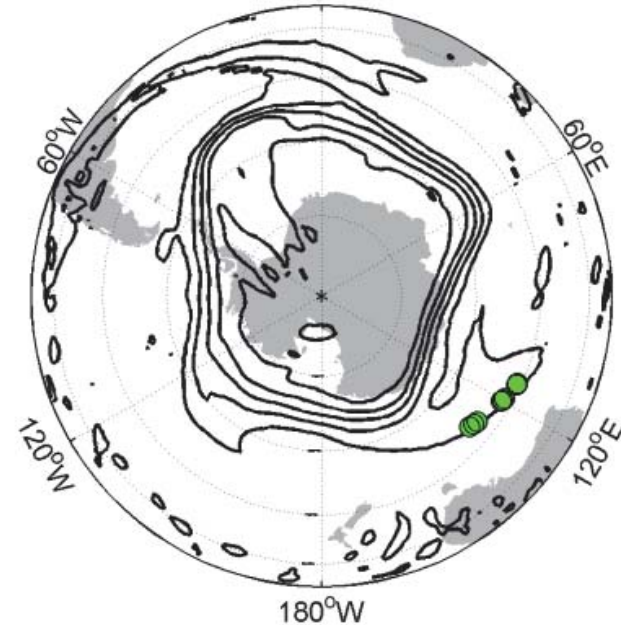
Potential vorticity fields

PV 17 Sep₀ 2005 00UTC



The balloons selected are located in the interior flank of the jet

PV 15 Oct₀ 2005 06UTC

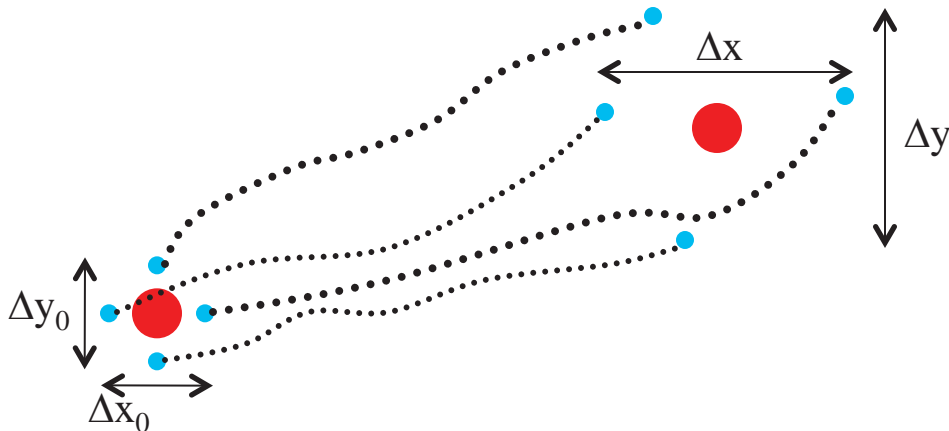


The balloons are expelled out of the vortex after ≈ 1 month



Finite-Time Lyapunov Exponents

FTLE are used to identify areas with maximum expansion rates between fluid particles initially close to each other.



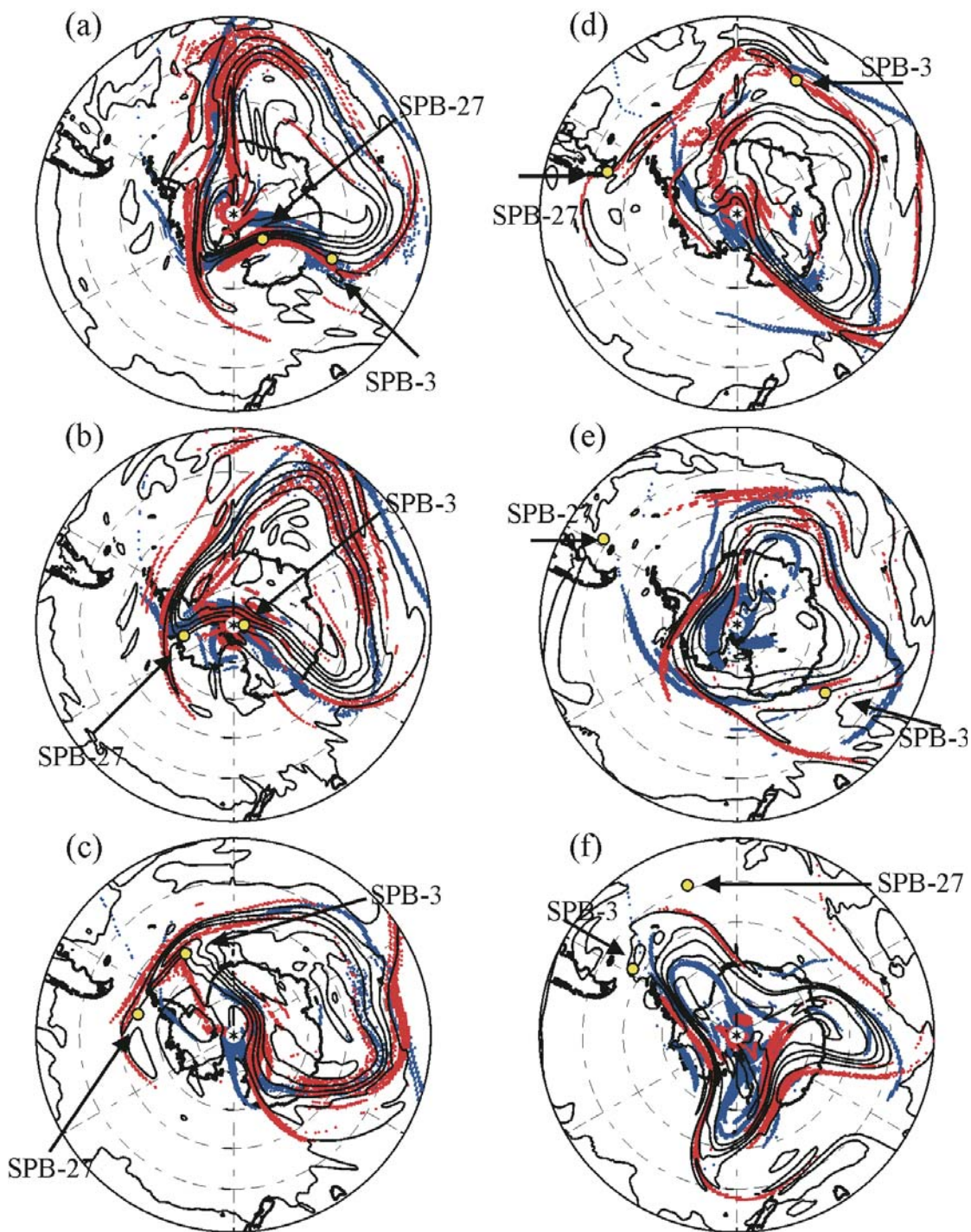
$$G = \begin{bmatrix} \frac{\Delta x}{\Delta x_0} & \frac{\Delta x}{\Delta y_0} \\ \frac{\Delta y}{\Delta x_0} & \frac{\Delta y}{\Delta y_0} \end{bmatrix}$$

Maximum eigenvalue of G^*G

$$LE = \frac{1}{\Delta t} \ln (\sqrt{\lambda_{max}})$$

Each particle is advected, forward and backward in time, during 4 days ($\Delta t = 4\text{days}$) with a 2nd-order Runge-Kutta-Hein algorithm integrating horizontal velocity field data at 2 different isopycnic levels (GEOS-5 reanalysis with VORCORE balloons).

Time step = 3h



Contours: potential vorticity maps on the isopycnic level $\rho = 0.0916 \text{ kg m}^{-3}$.

Contour interval 8 PVU (1 PVU = $10^{-6} \text{ m}^2 \text{ s}^{-1} \text{ kg}^{-1} \text{ K}$).
 Shaded Central initial location of set of particles with FTLE values above 0.033 h^{-1} for the forward (blue) and backward (red) integrations at the same isopycnic level.
 Data from ECMWF

Yellow dots

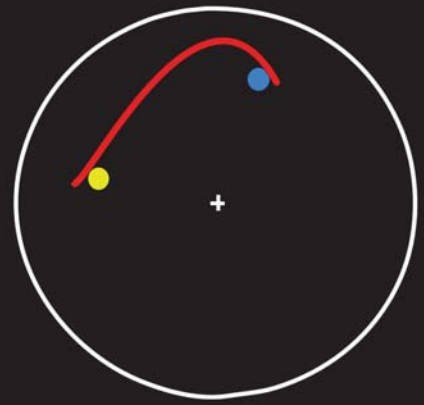
Location of SPB-3 and SPB-27

- (a) 21 November at 06 UTC,
- (b) 22 November at 06 UTC,
- (c) 23 November at 18 UTC,
- (d) 25 November at 18 UTC,
- (e) 28 November at 09 UTC,
- (f) 2 December at 00 UTC

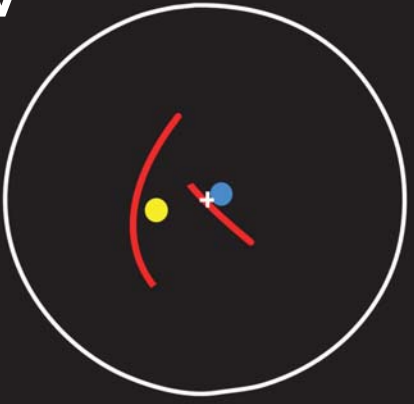
21 Nov



25 Nov



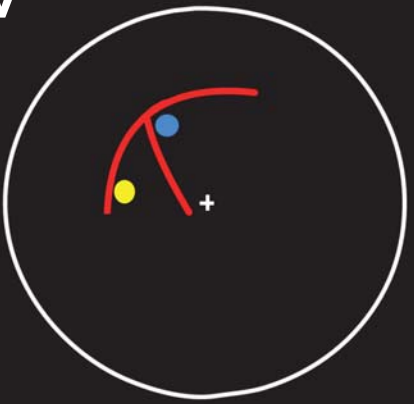
22 Nov



28 Nov

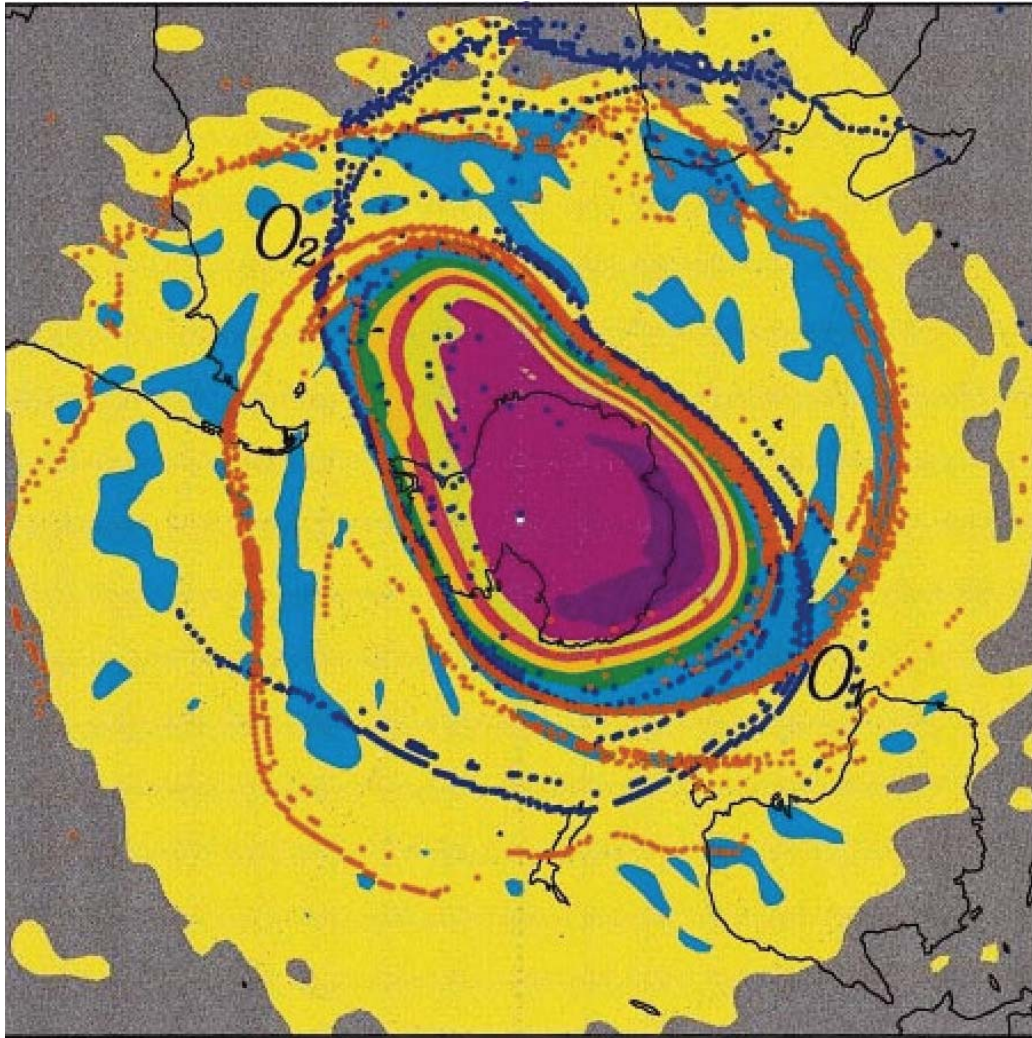


23 Nov



2 Dec





PV on 500
Data : ECMWF

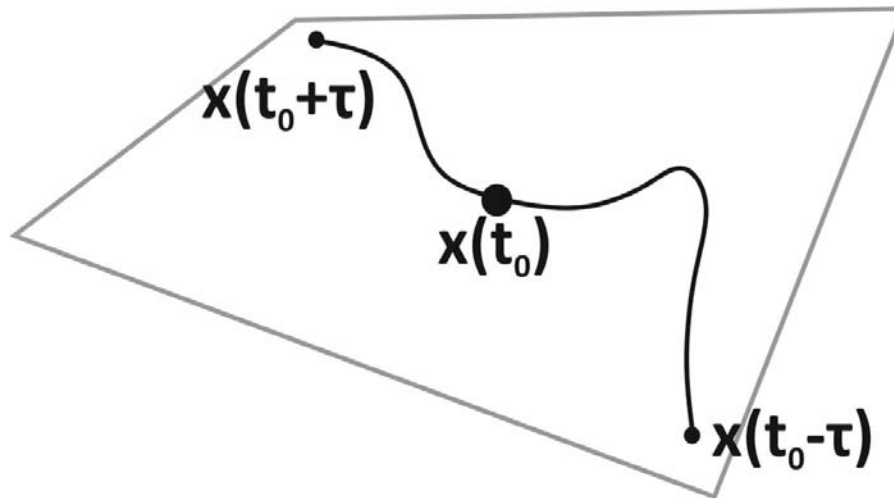
Red dots: Unstable manifold

Blue dots: Stable manifold

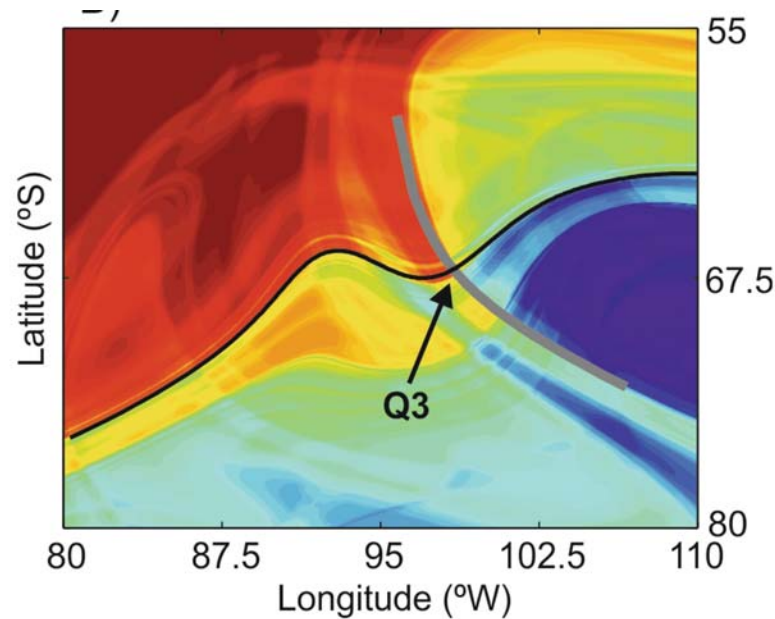
The Lagrangian descriptor M

M measures the length of trajectories over a range of times $[t_0 - \tau, t_0 + \tau]$ passing through $(x(t_0), y(t_0), z(t_0))$.

$$M_\tau(x_0, y_0, z_0, t_0) = \int_{t_0 - \tau}^{t_0 + \tau} dt \sqrt{\left(\frac{dx(t)}{dt}\right)^2 + \left(\frac{dy(t)}{dt}\right)^2 + \left(\frac{dz(t)}{dt}\right)^2}$$



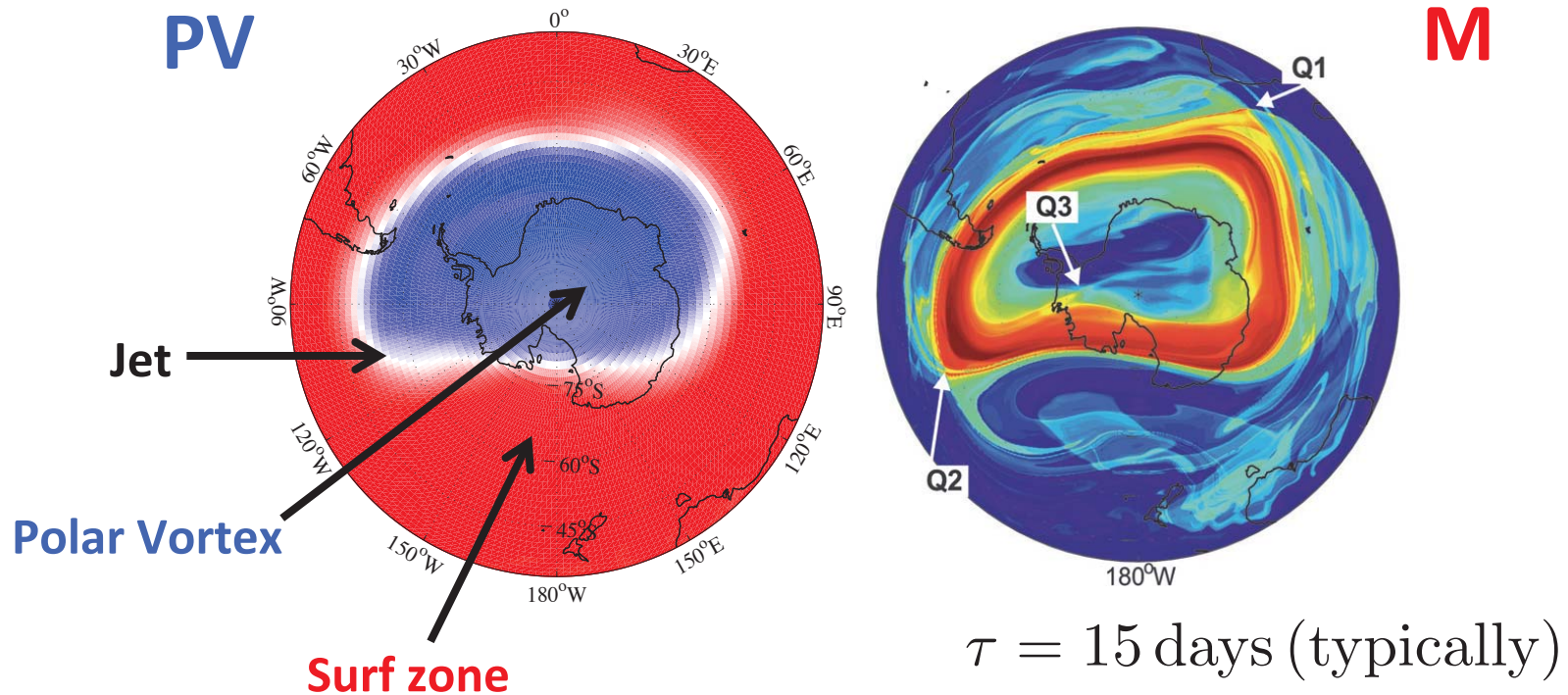
Crosshatched patterns in M



Hyperbolic points (Q_3) are at the intersection of stable and unstable manifolds (solid lines).

In unsteady flows they define hyperbolic trajectories.

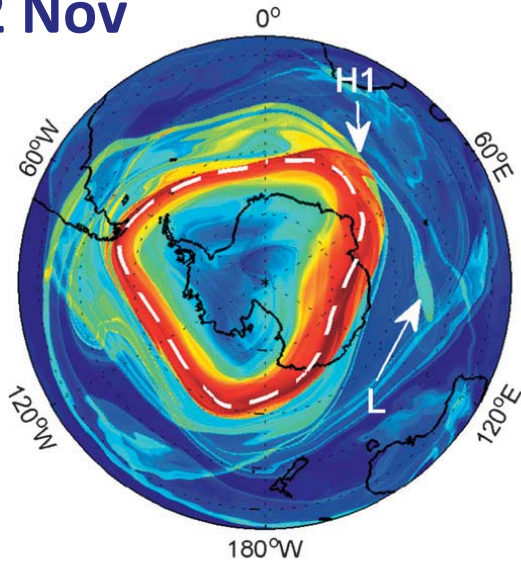
ERA-INTERIM data: OCTOBER 8, 2005, 475K



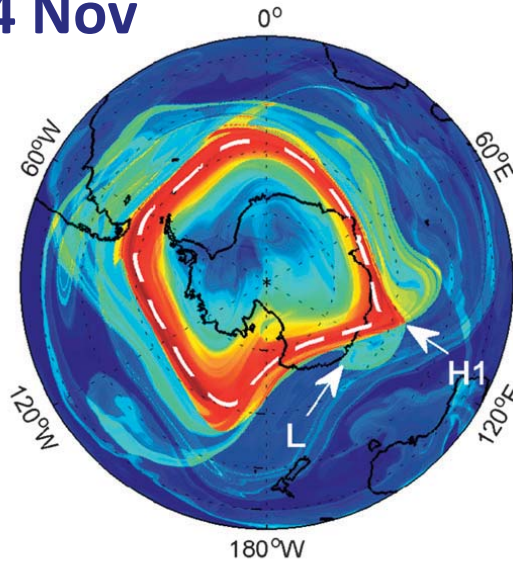
- De La Camara et al. (2012, 2013) found **hyperbolic points** both **outside** (Q1, Q2) and **inside** (Q3) the polar vortex.
- They suggested that intersection of lobes from outside and inside the vortex provides a route of transport across the vortex edge.

M from ERA-INTERIM data: NOVEMBER 2005, 475K

2 Nov

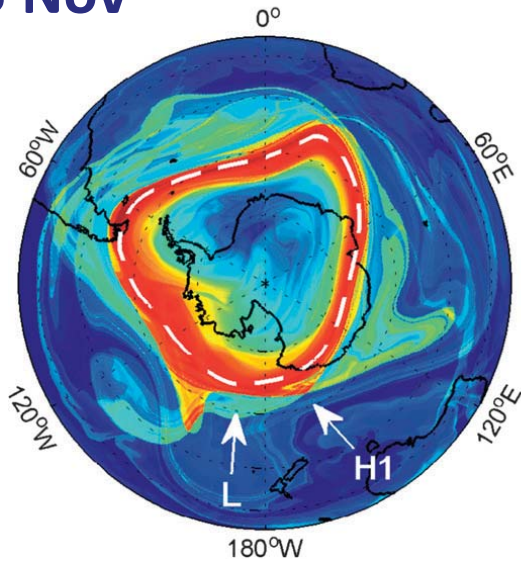


4 Nov

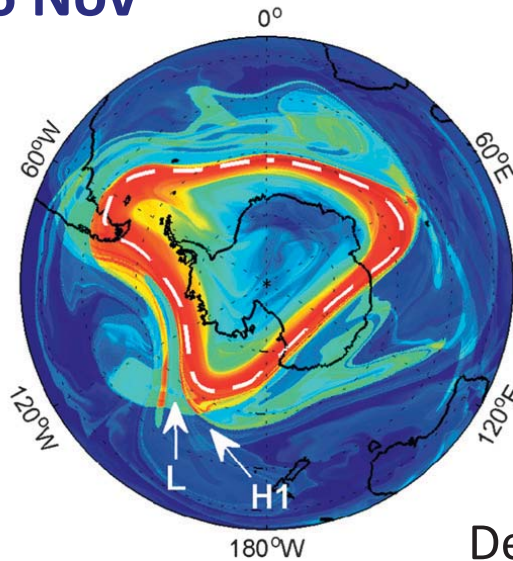


H: Hyperbolic Point

5 Nov



6 Nov

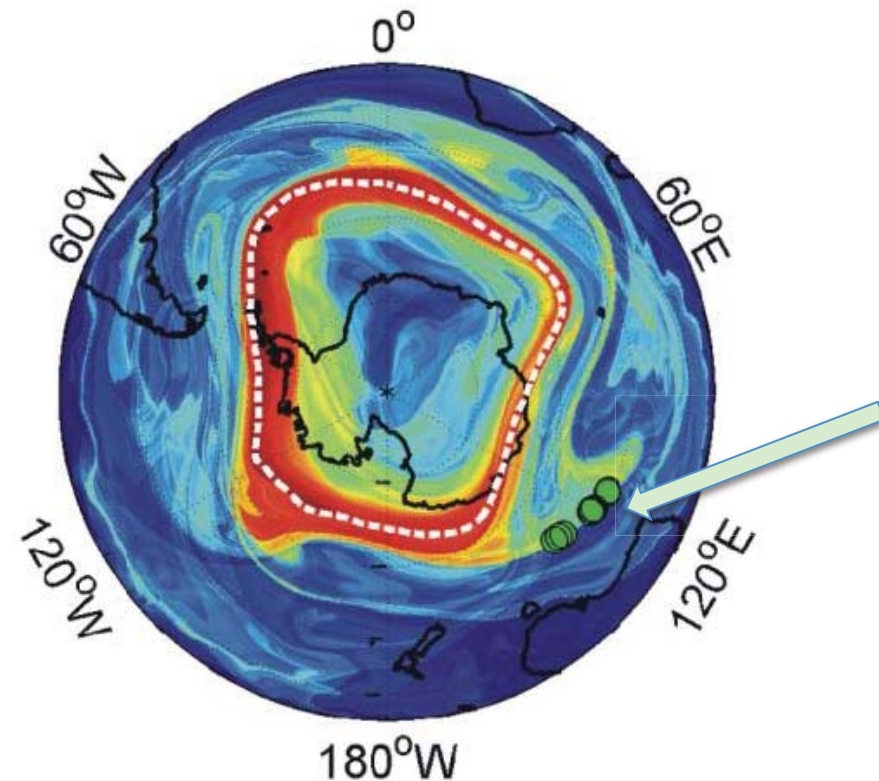
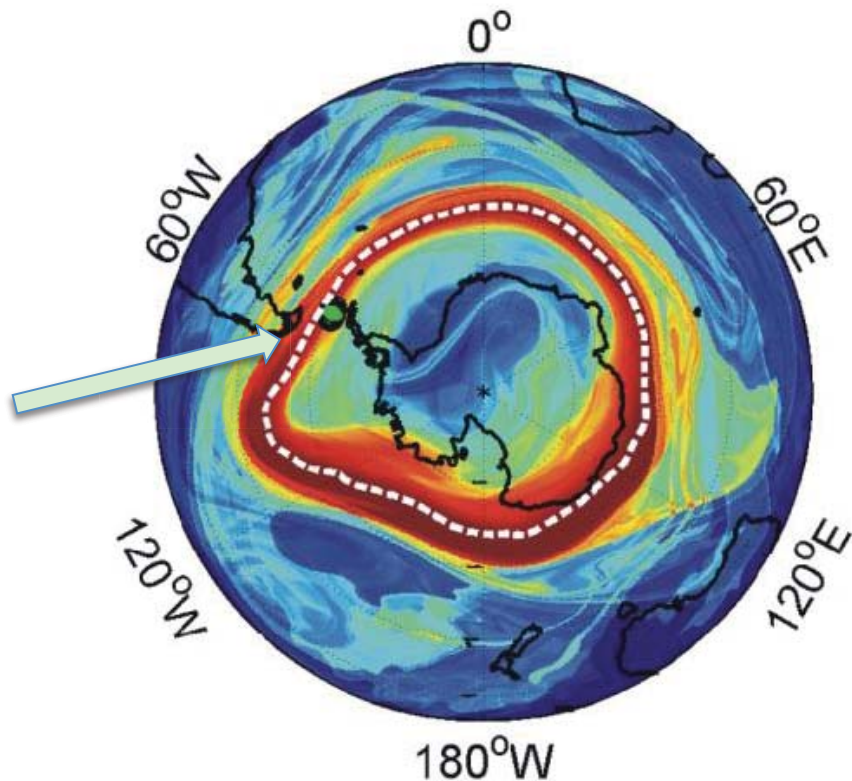


L: Lobe

Fluid parcels leaving the Vortex

M: 17/Sep/05; $\tau=10$; Forward

M: 15/Oct/05; $\tau=10$; Backward

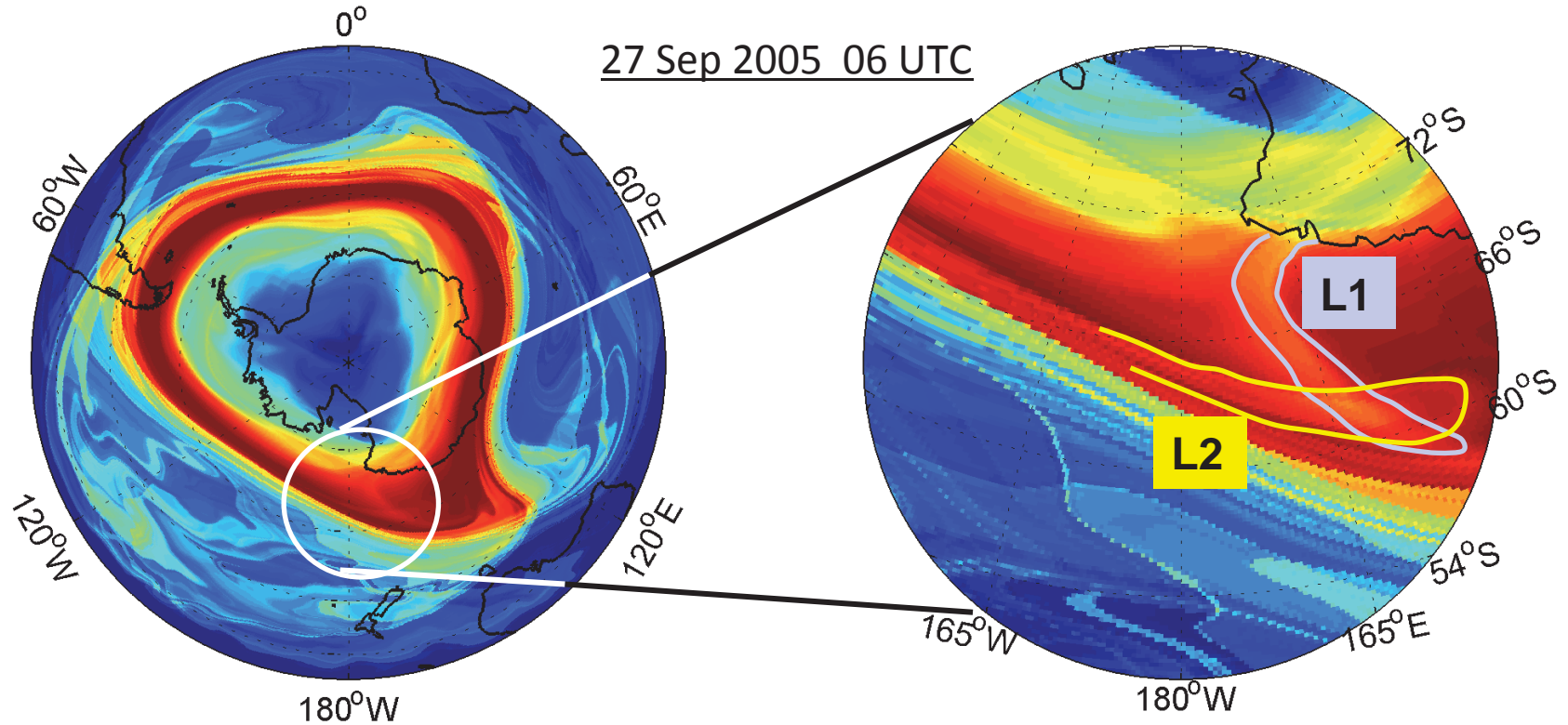


475K; Dotted line indicates max PV gradient

Transport **ACROSS** the stratospheric polar vortex

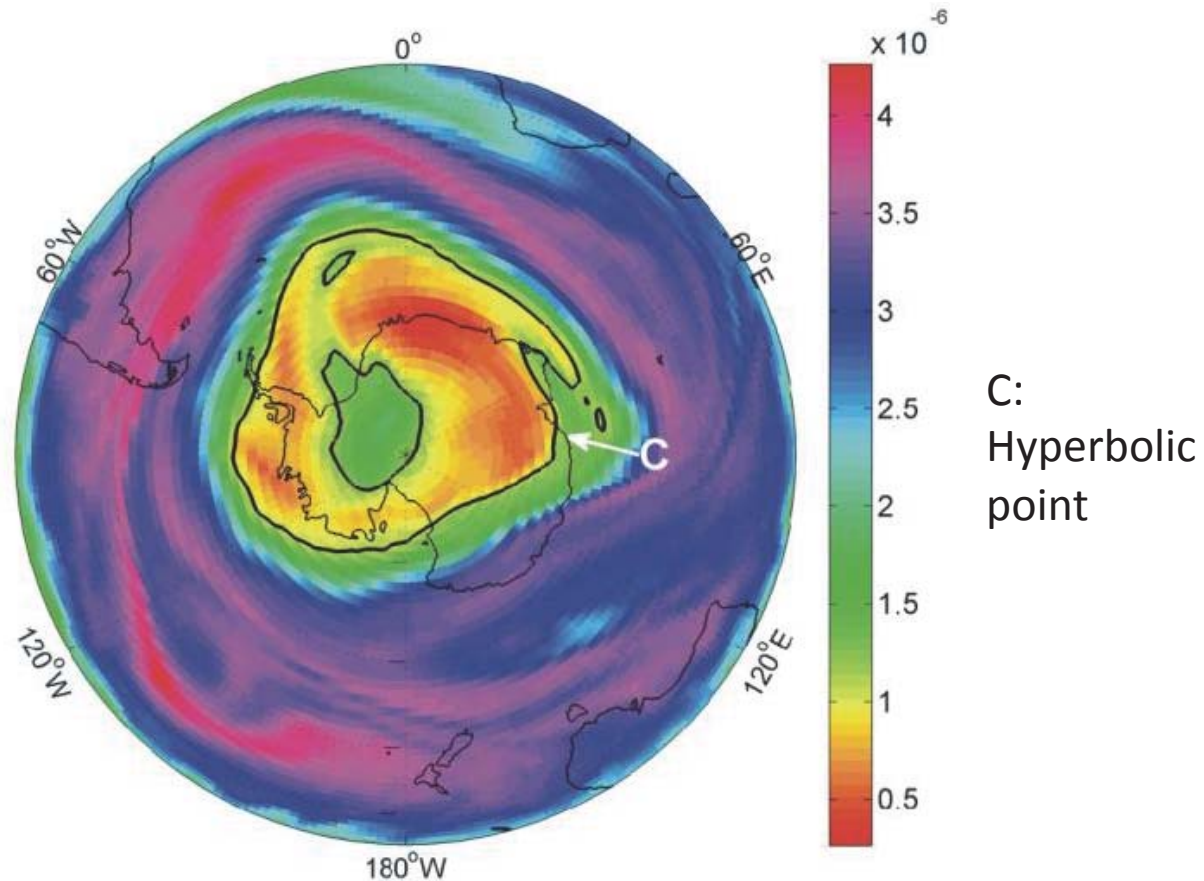
Descriptor M, $\tau=15$ days

27 Sep 2005 06 UTC



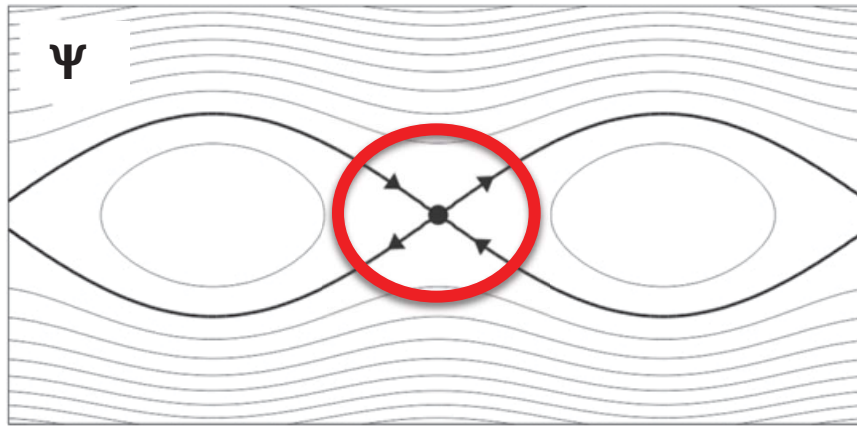
Intersections of lobes in the edge region that lead to effective mass transport across the vortex edge are identified.

Signature of lobes in Ozone field

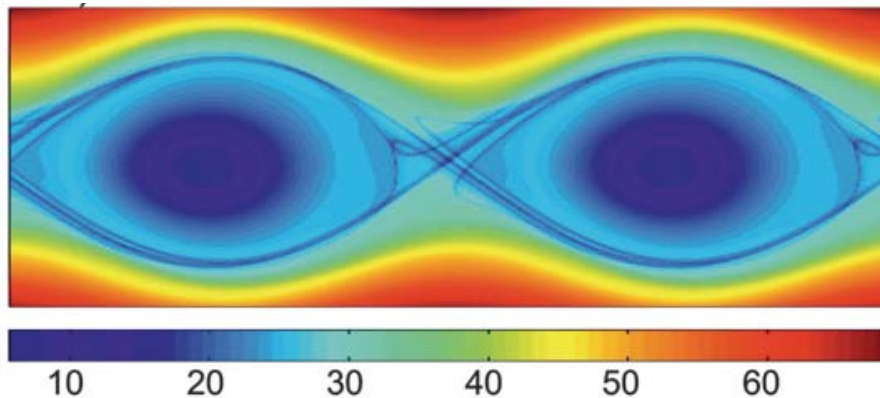


Ozone mass mixing ratio at 475K for 17 October 2005

“CAT EYES” AND LAGRANGIAN DESCRIPTOR “M”

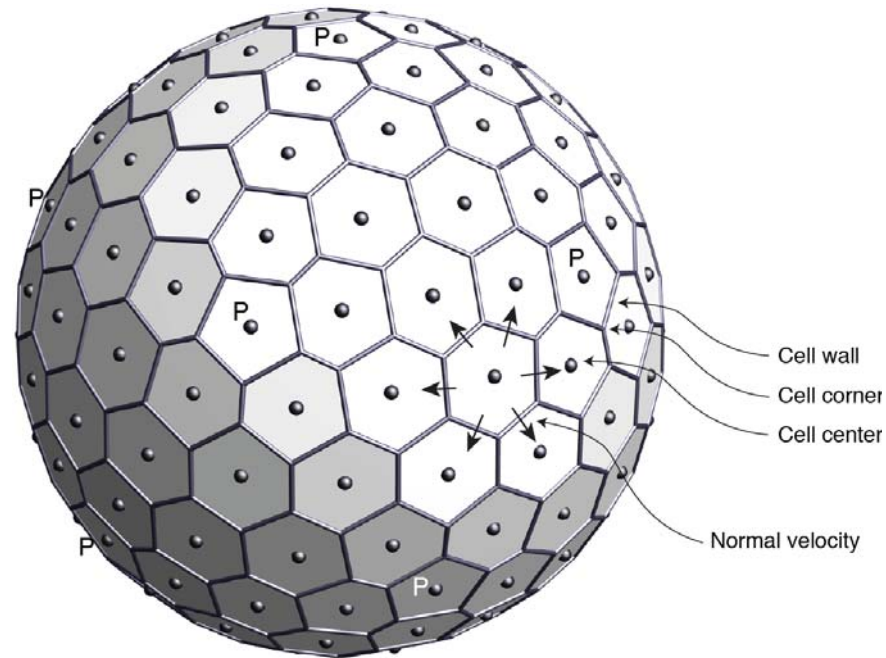


$$M_\tau(x_0, y_0, z_0, t_0) = \int_{t_0 - \tau}^{t_0 + \tau} \sqrt{\left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2 + \left(\frac{dz}{dt}\right)^2} dt$$



M shows the Lagrangian skeleton of the flow field (Madrid and Mancho; Chaos 2009).

Methodology: We perform integrations with a one-layer, shallow-water model on the sphere from realistic, zonally independent initial conditions

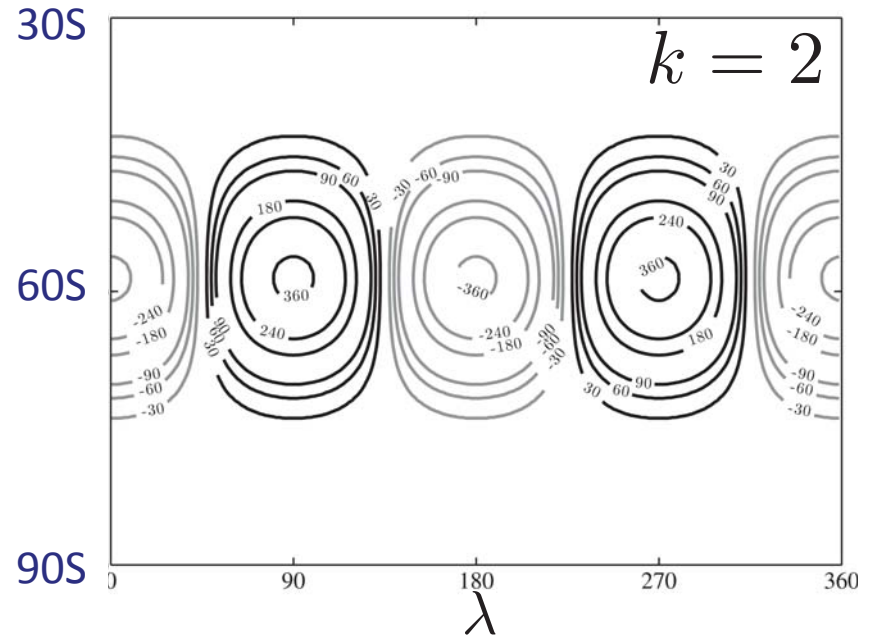
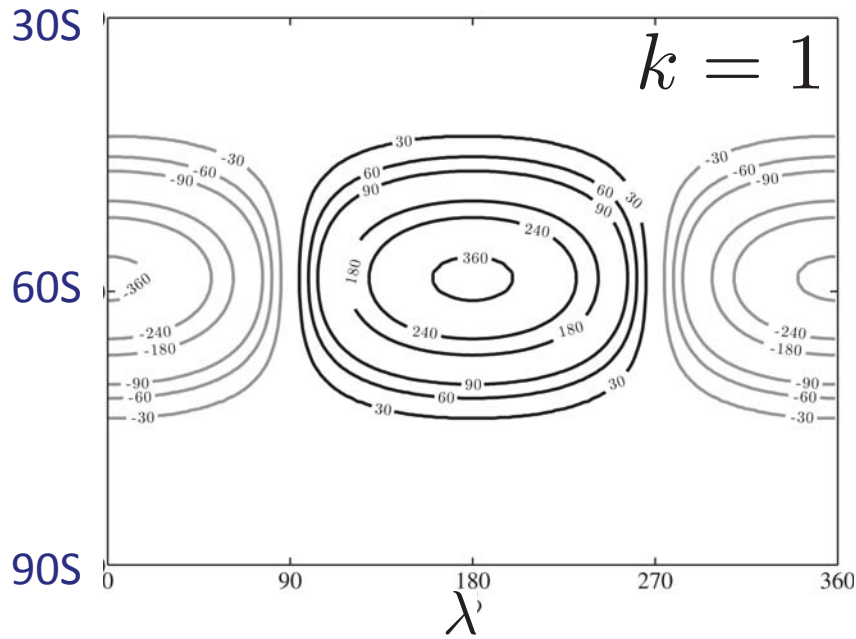


40962 grid cells yielding approx 120km grid distance

Model description in Heikes, Randall and Konor (Mon. Wea. Rev., 2013).

For display, model results are interpolated to a long-lat grid.

Model's bottom topography



$$h_b = -B \left[\frac{\cos(\phi)}{\cos(\phi_0)} \right] \exp \left[- \left(\frac{\phi - \phi_0}{\Delta\phi} \right)^2 \right] \left[1 - \exp \left(- \frac{t}{t_S} \right) \right] \cos [k\lambda - \omega t]$$

Nearly Gaussian Profile

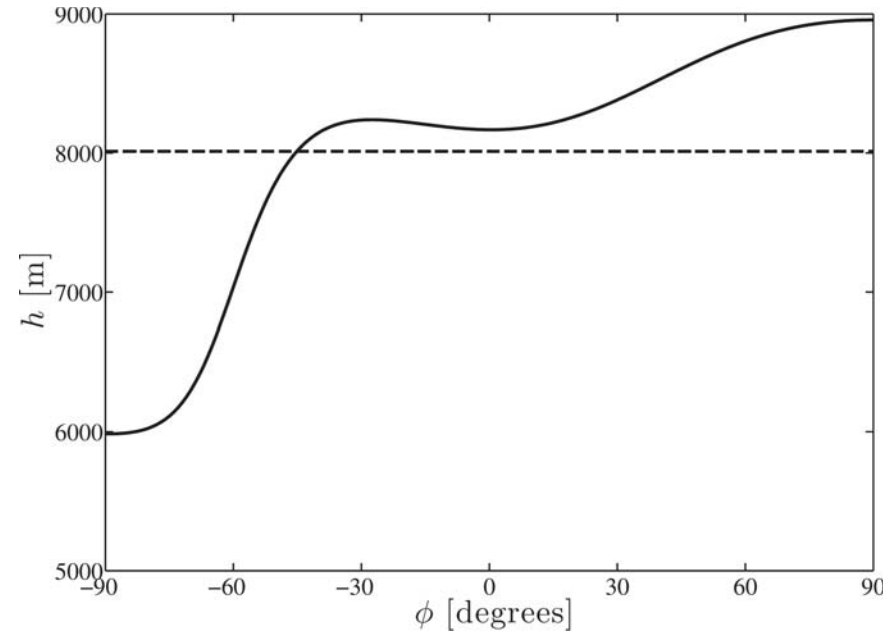
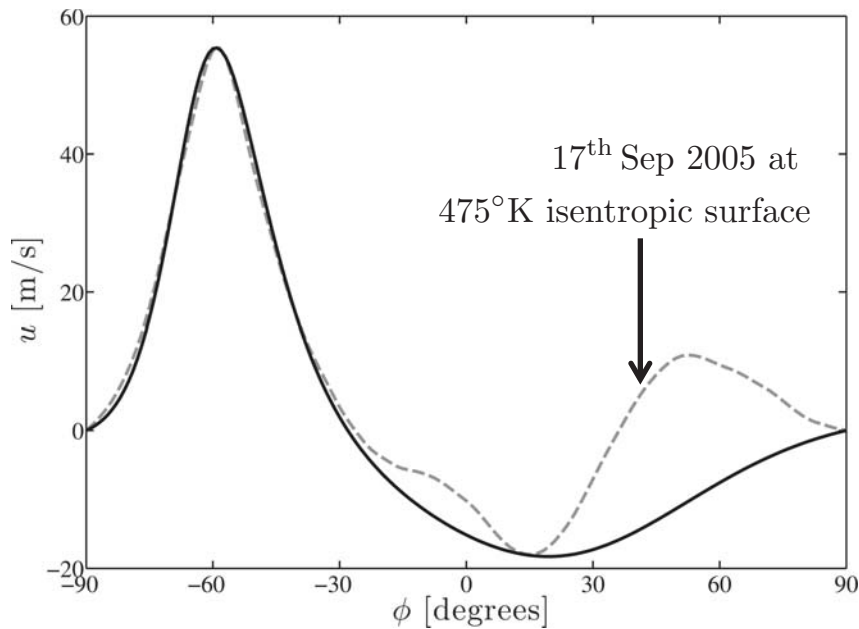
Grows in Time

Can Travel

$$B = 392\text{m}, \phi_0 = -\pi/3, \Delta\phi = \pi/18, t_S = 10\text{days}$$

(Ngan and Shepherd, 1999)

Initial velocity and surface height profiles

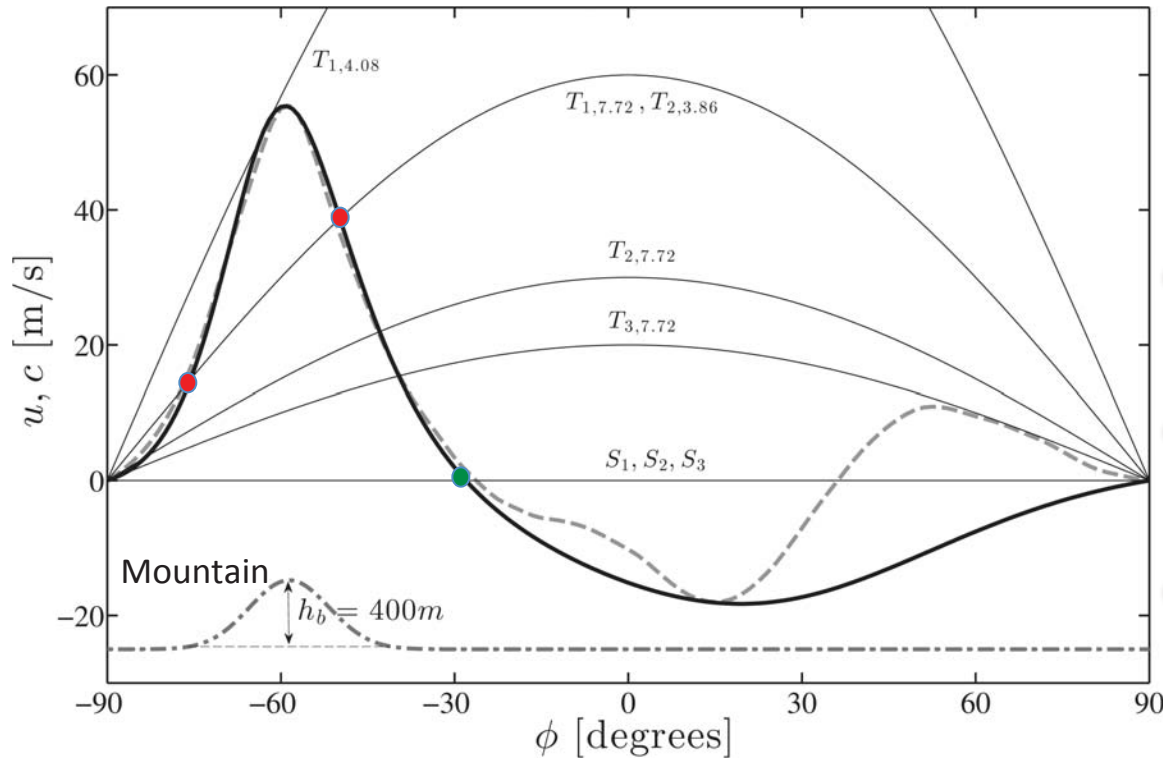


$$u(\phi) = 118 \cos(\phi) \operatorname{sech}\left(\frac{\phi + 1.09}{0.185}\right) - 20 \cos(\phi) \operatorname{sech}\left(\frac{0.5 - \phi}{0.7}\right).$$

$$h(\phi) = 5985 - \int^{\phi} \frac{Ru(\phi')}{g} \left[f + \frac{\tan(\phi')}{R} u(\phi') \right] d\phi'$$

Mean height of shallow water is 8km.

Selected Experiments



| Case | k | ω (day ⁻¹) | T (day) |
|--------------|-----|----------------------------------|--------------|
| S_1 | 1 | 0 | ∞ |
| $T_{1,7.72}$ | 1 | 0.81 | 7.72 |
| $T_{1,4.08}$ | 1 | 1.53 | 4.08 |
| S_2 | 2 | 0 | ∞ |
| $T_{2,7.72}$ | 2 | 0.81 | 7.72 |
| $T_{2,3.86}$ | 2 | 1.63 | 3.86 |
| S_3 | 3 | 0 | ∞ |
| $T_{3,7.72}$ | 3 | 0.81 | 7.72 |

Phase-speed

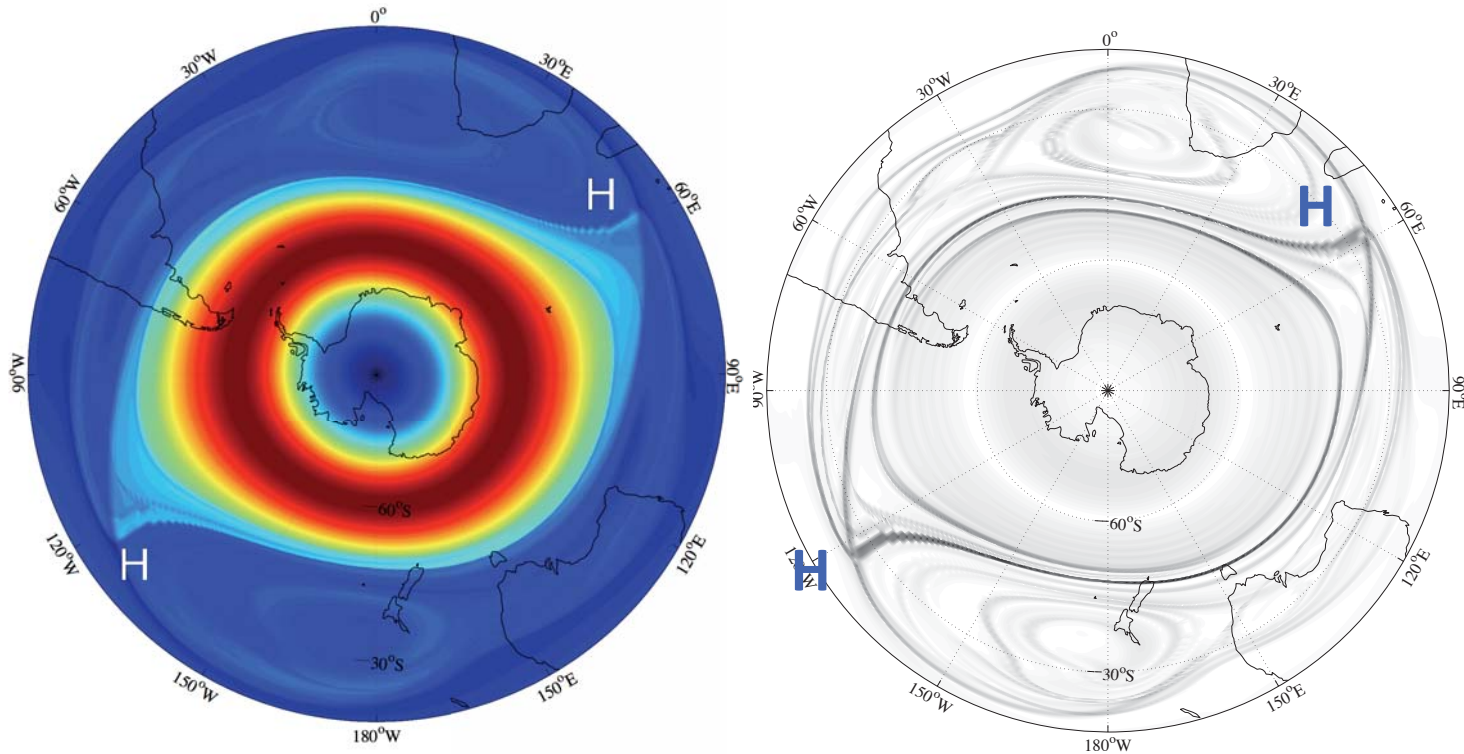
$$c = \frac{\omega R \cos(\phi)}{k}$$

Forcings at different zonal wave numbers and speeds consistent with one or two critical layers

M and $|\nabla(M)|$ plots

(Contours of $|\nabla(M)|$ align with the manifolds)

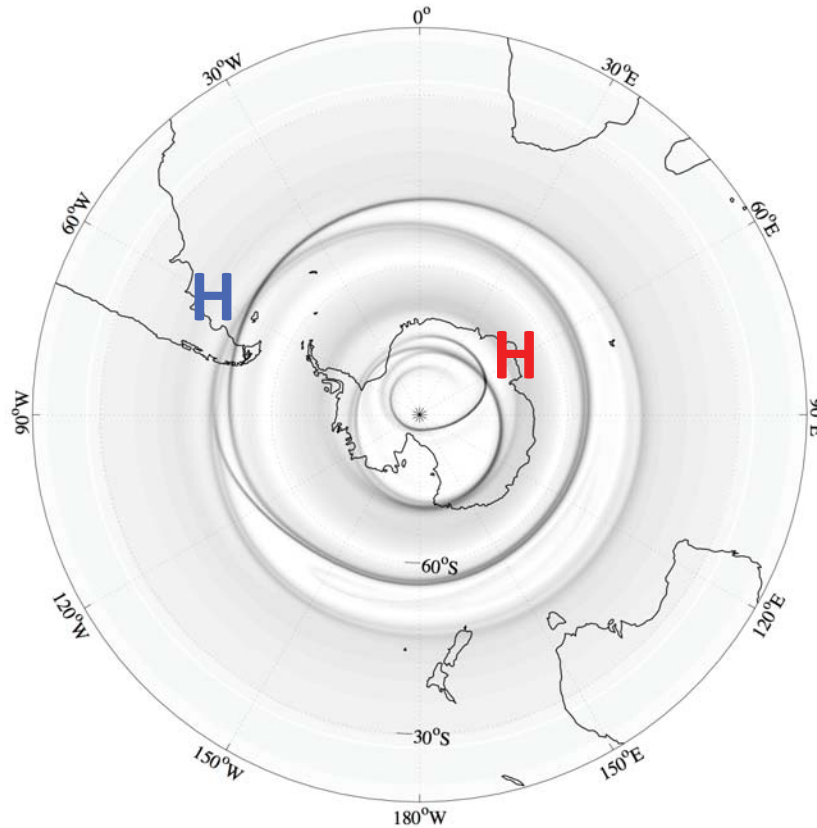
S_2 : Stationary wavenumber 2



Note the two hyperbolic points at the latitude where $u=0$

M and $|\nabla(M)|$ plots

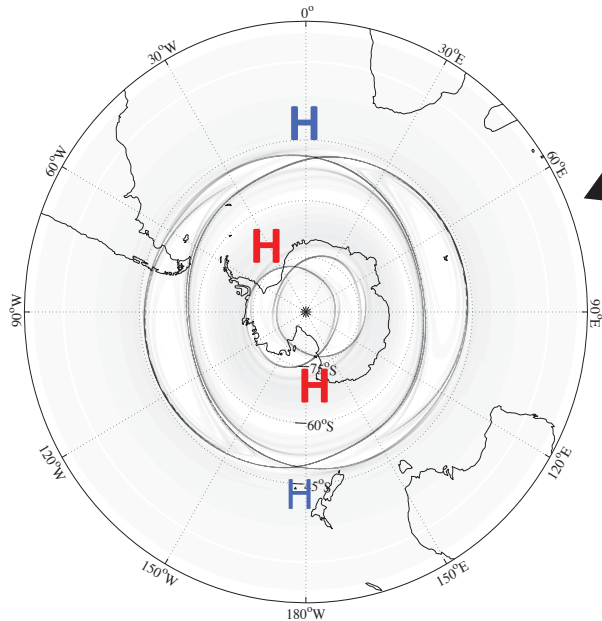
$T_{1,7.72}$: Traveling wavenumber 1 with period 7.72 d



Note the two hyperbolic points, one outside and the other inside the vortex, at the critical latitudes.

M and $|\nabla(M)|$ plots

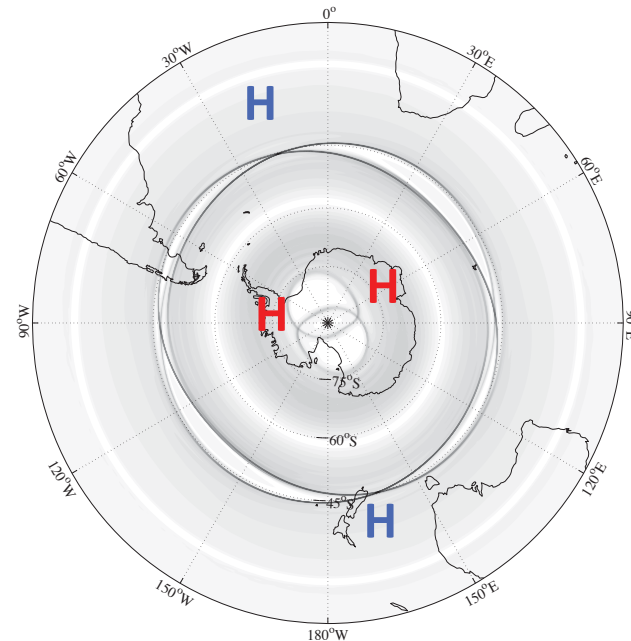
$T_{2,3.86}$: Traveling wavenumber 2 with period 3.86 d



Four hyperbolic points, two outside and two inside the vortex, at the critical latitudes

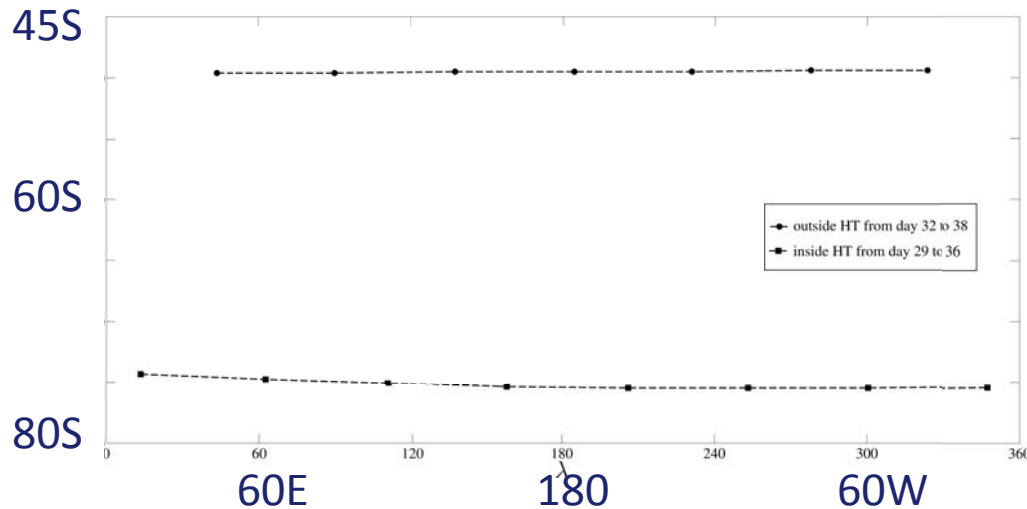
$T_{2,3.86}$, except for mountain centered at 45S (and completely outside SPV)

Still four hyperbolic points, two outside and two inside the vortex at the critical latitudes, even though forcing is outside the vortex

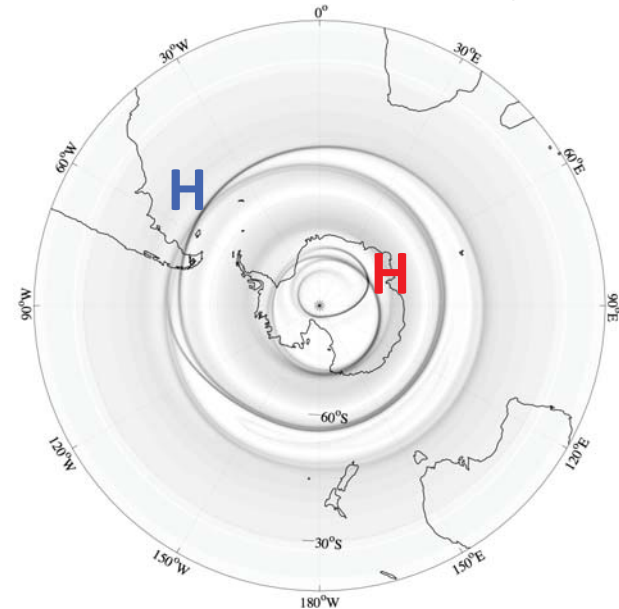


$T_{1,7.72}$: Traveling wavenumber 1 with period 7.72 d

Hyperbolic Trajectories
(points are 1 day apart)

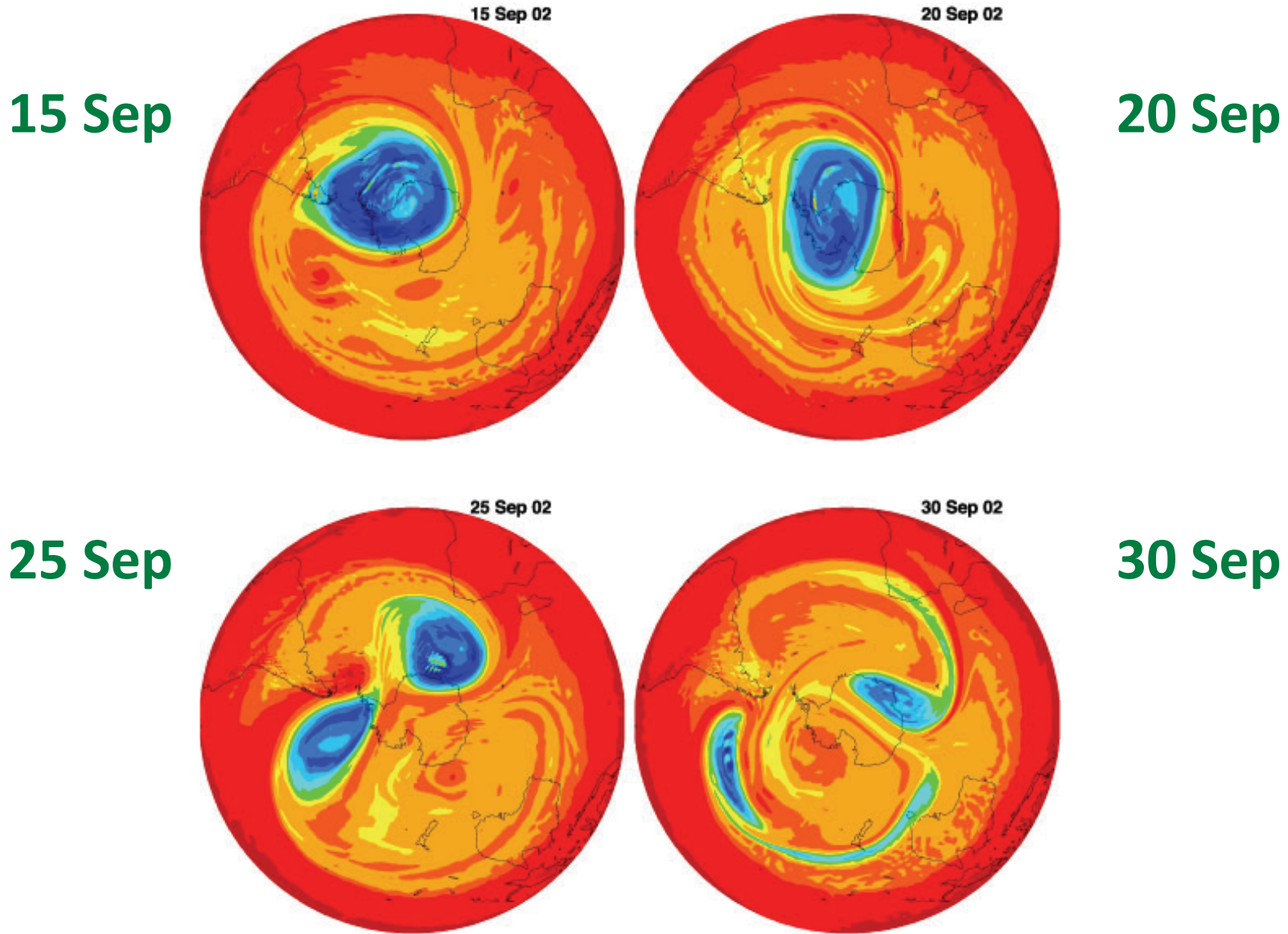


Plot of $|\nabla(M)|$ on day 25



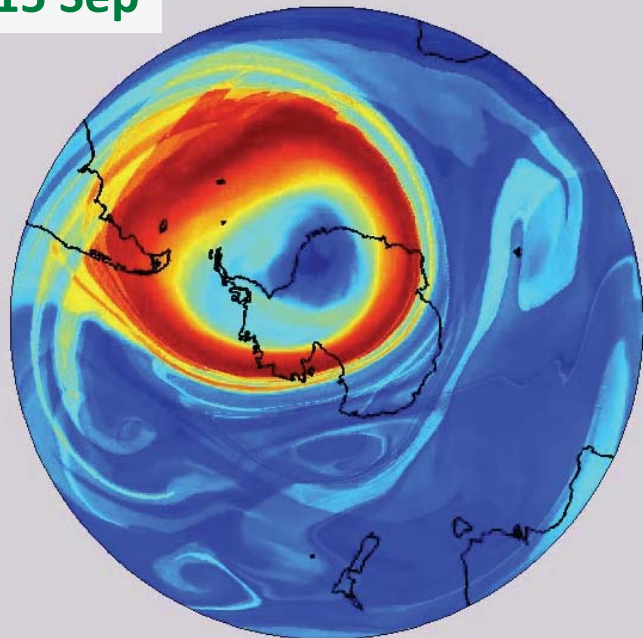
Hyperbolic trajectories approximately coincide with
the wave's critical latitudes for several days
HTs provide an indirect measure of the phase speed
of the breaking waves.

The unique event in 2002: PV at 850K



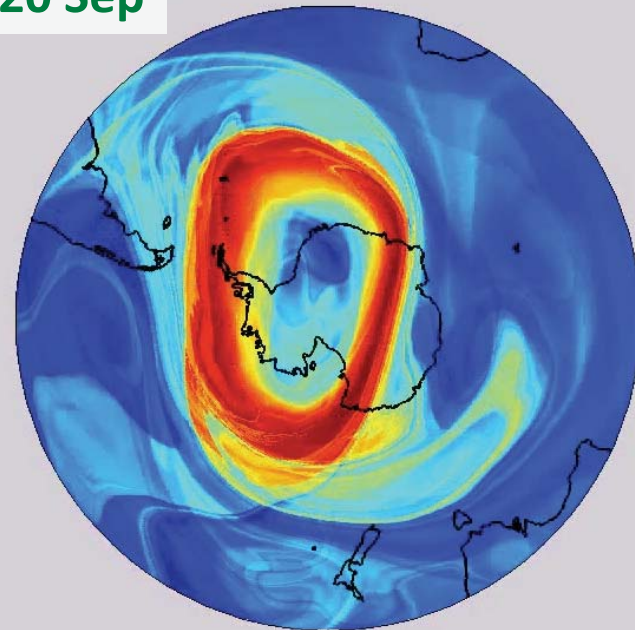
M 15 Sep

for tau = 10 at 15/9/2002 12:00:00 UTC (850 K)



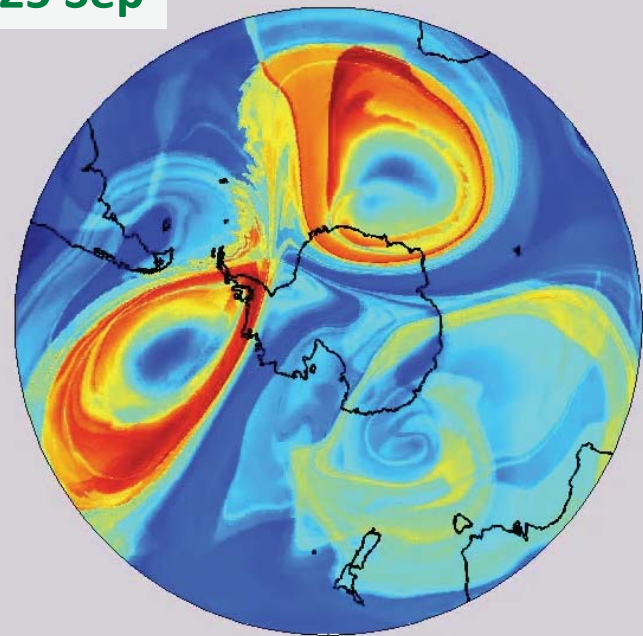
M 20 Sep

for tau = 10 at 20/9/2002 12:00:00 UTC (850 K)



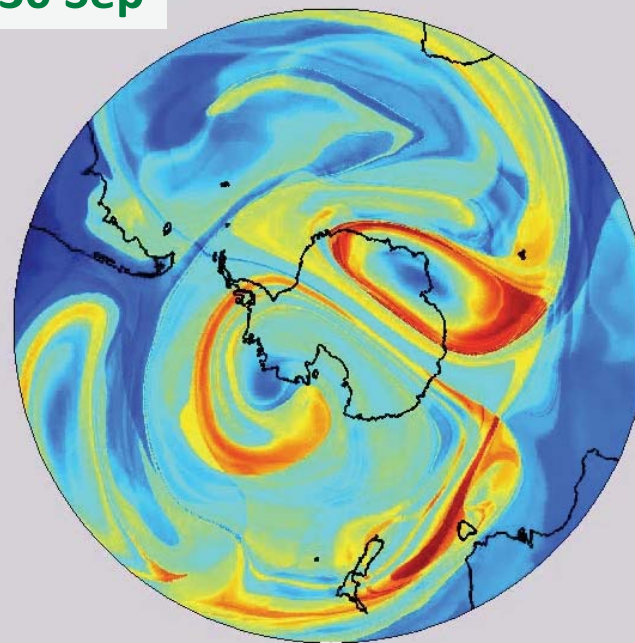
M 25 Sep

for tau = 10 at 25/9/2002 12:00:00 UTC (850 K)



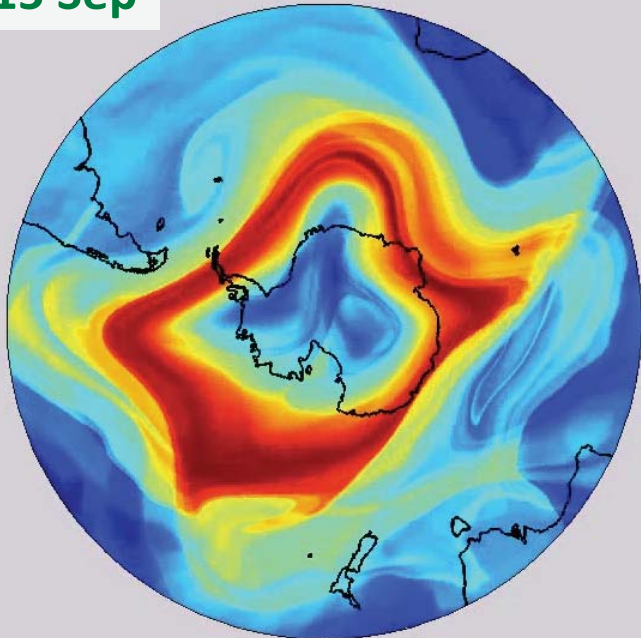
M 30 Sep

for tau = 10 at 30/9/2002 12:00:00 UTC (850 K)



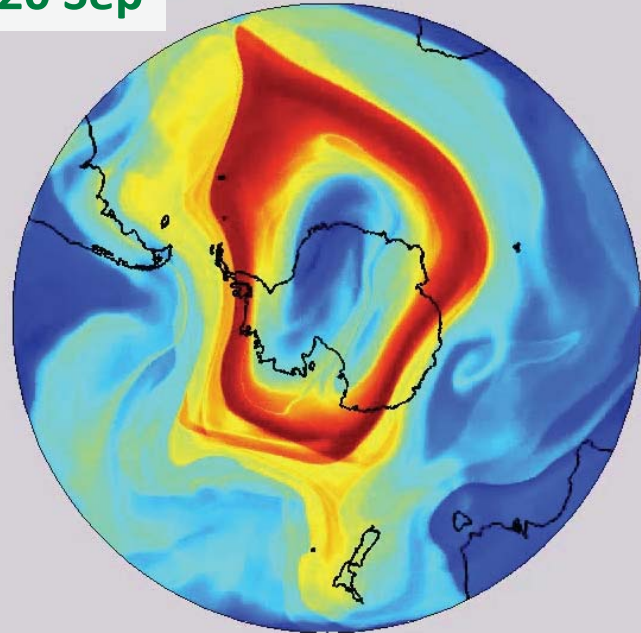
M 15 Sep

for $\tau = 10$ at 15/9/2002 12:00:00 UTC (430 K)



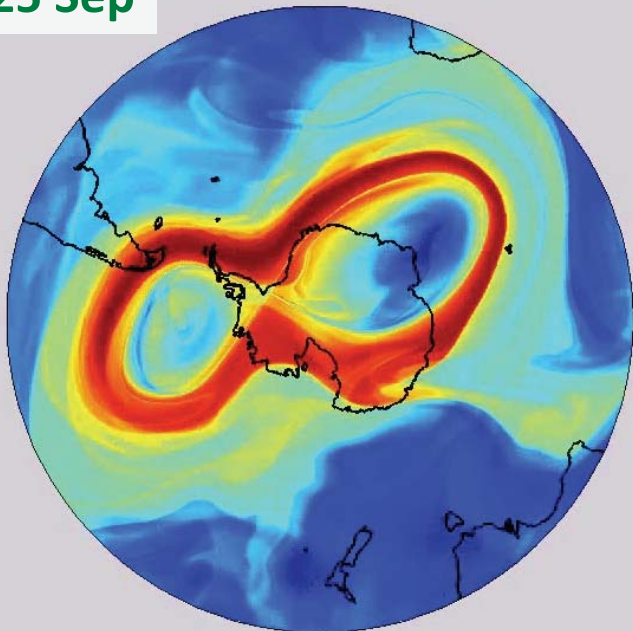
M 20 Sep

for $\tau = 10$ at 20/9/2002 12:00:00 UTC (430 K)



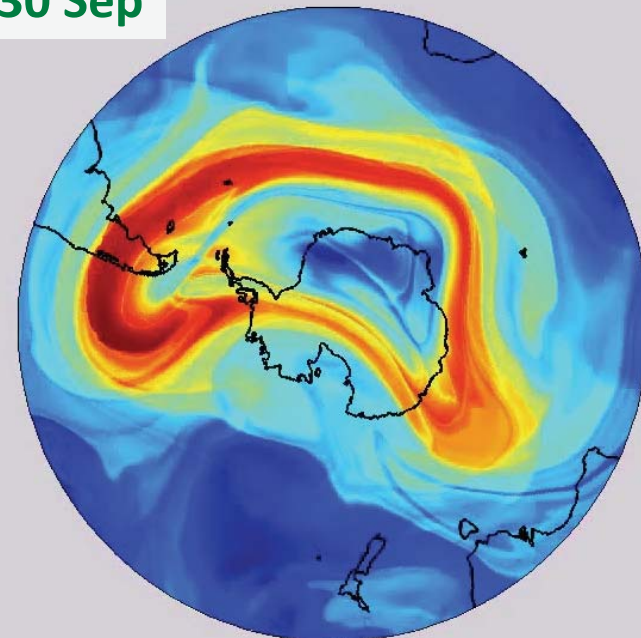
M 25 Sep

for $\tau = 10$ at 25/9/2002 12:00:00 UTC (430 K)

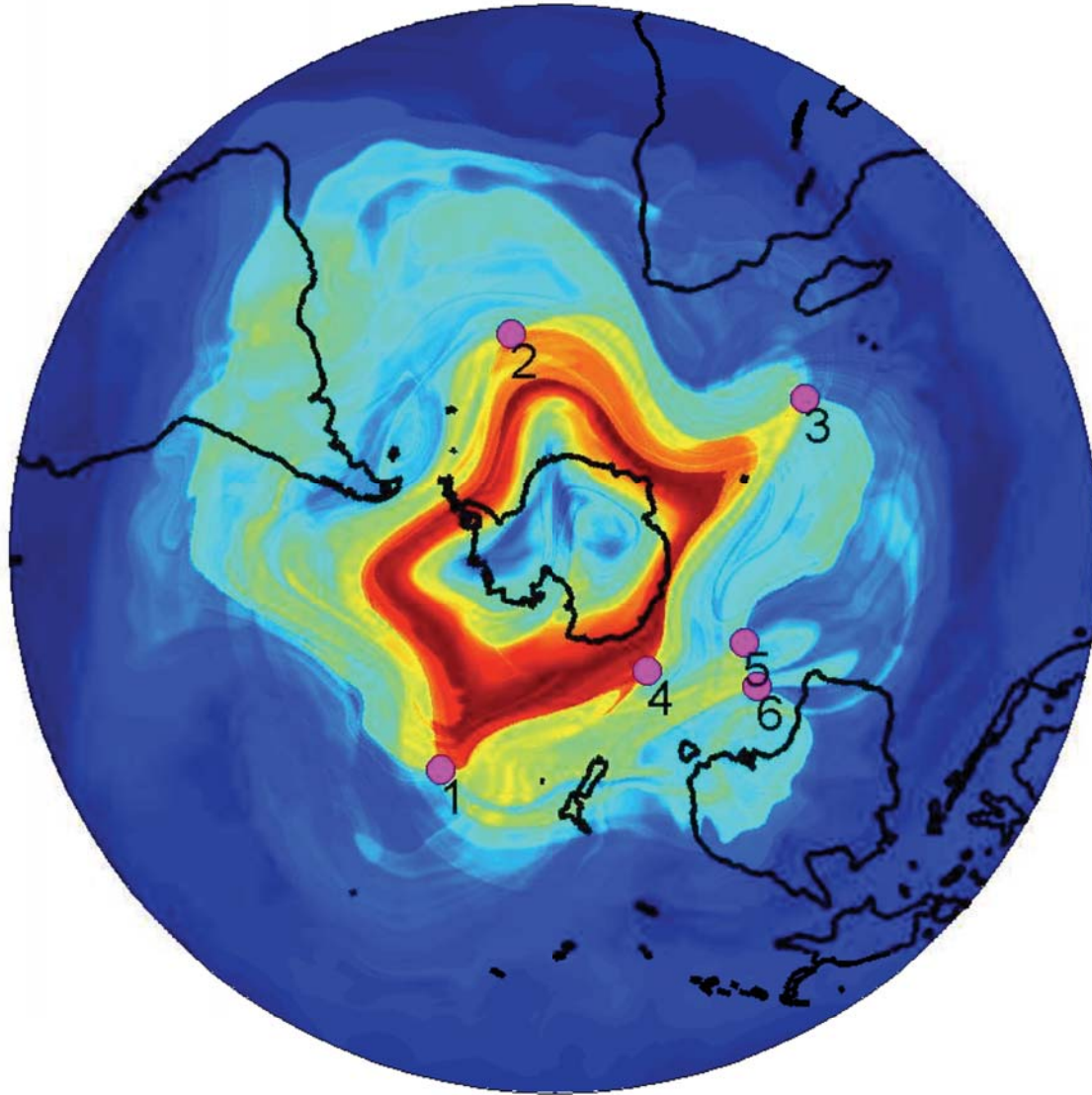


M 30 Sep

for $\tau = 10$ at 30/9/2002 12:00:00 UTC (430 K)



Hyperbolic points on 15 Sep 2002, 430K



Synthesis

- The most dramatic changes in the Antarctic stratosphere occur in Spring. Development of the “Antarctic Ozone Hole” motivated studies on transports in the stratosphere.
- We do not recognize “bipartidismo” (M or LE) for Lagrangian analysis: Using the *popular M, citizens united can* find hyperbolic point(s)/trajectories both **outside** and **inside** the southern Stratospheric Polar Vortex (SPV)
- Rossby Wave breaking (RWB) was identified hyperbolic points. These remain at the wave’s critical latitudes for several days.
- Current research focuses on the 3D structures of RWB and associated lobes.

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